

Two-dimensional turbulence above topography



a 1976 JFM paper by
F. P. Bretherton & D. H. Haidvogel
(as told by Cesar)

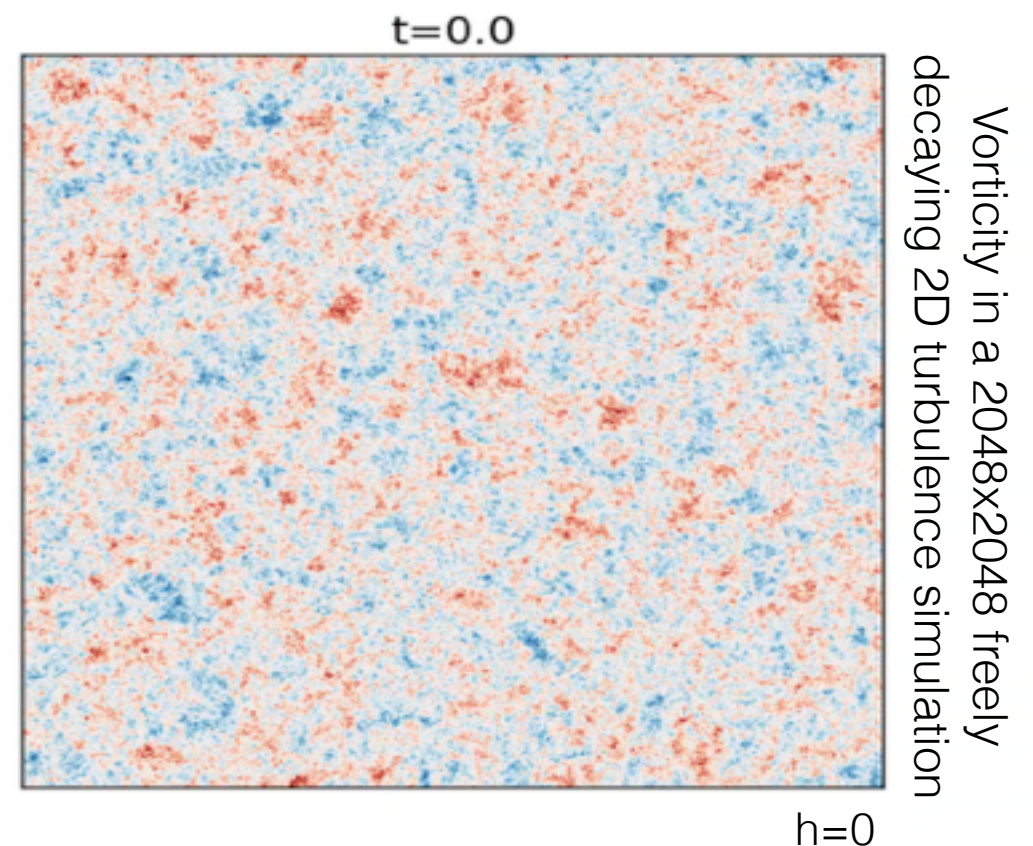


Context

MODE experiment

FFT (1965), 2D turbulence studies

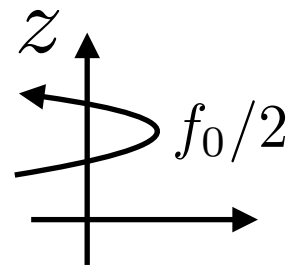
Contemporary studies: Rhines (1975, 1977),
Salmon (1978, 1980), McWilliams (1984) →



Inviscid 2D dynamics

(f-plane barotropic QG dynamics)

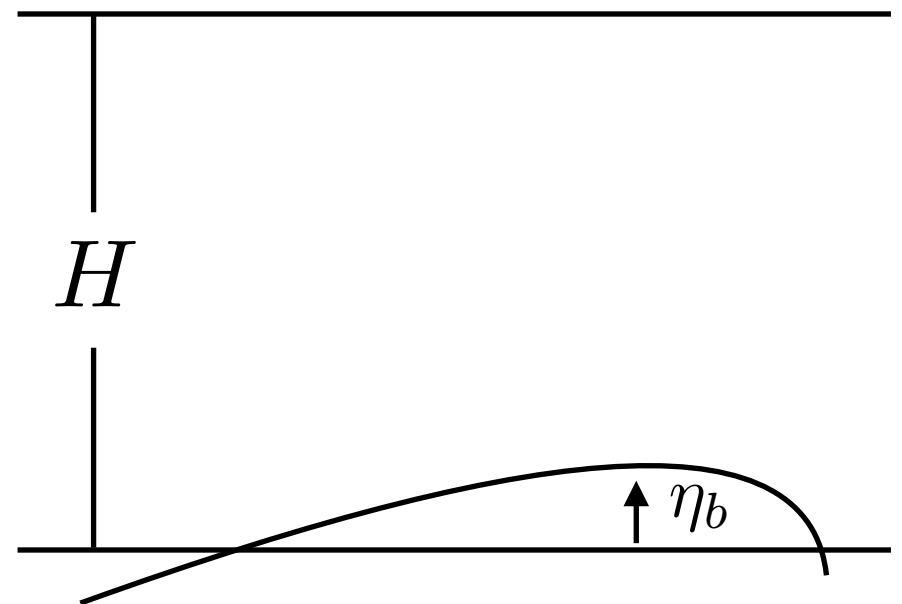
$$\frac{Dq}{Dt} = \partial_t q + \mathbf{J}(\psi, q) = 0$$



$$q = \nabla^2 \psi + h$$

$$u = -\psi_x \quad v = \psi_y$$

$$\mathbf{J}(f, g) = f_x g_y - f_y g_x$$

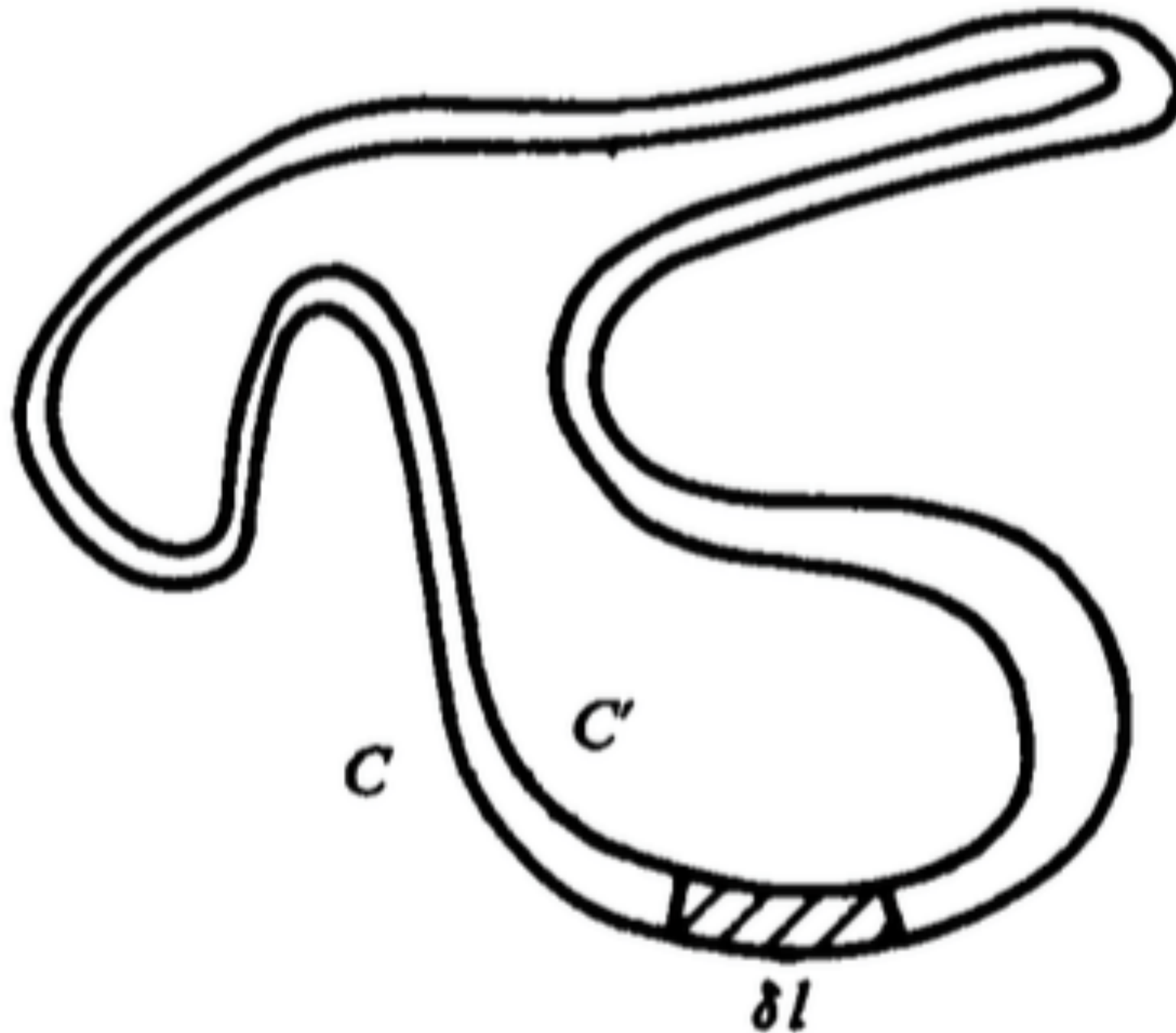


Note $h = \frac{f_0}{H} \eta_b$

The enstrophy cascade

2D turbulence

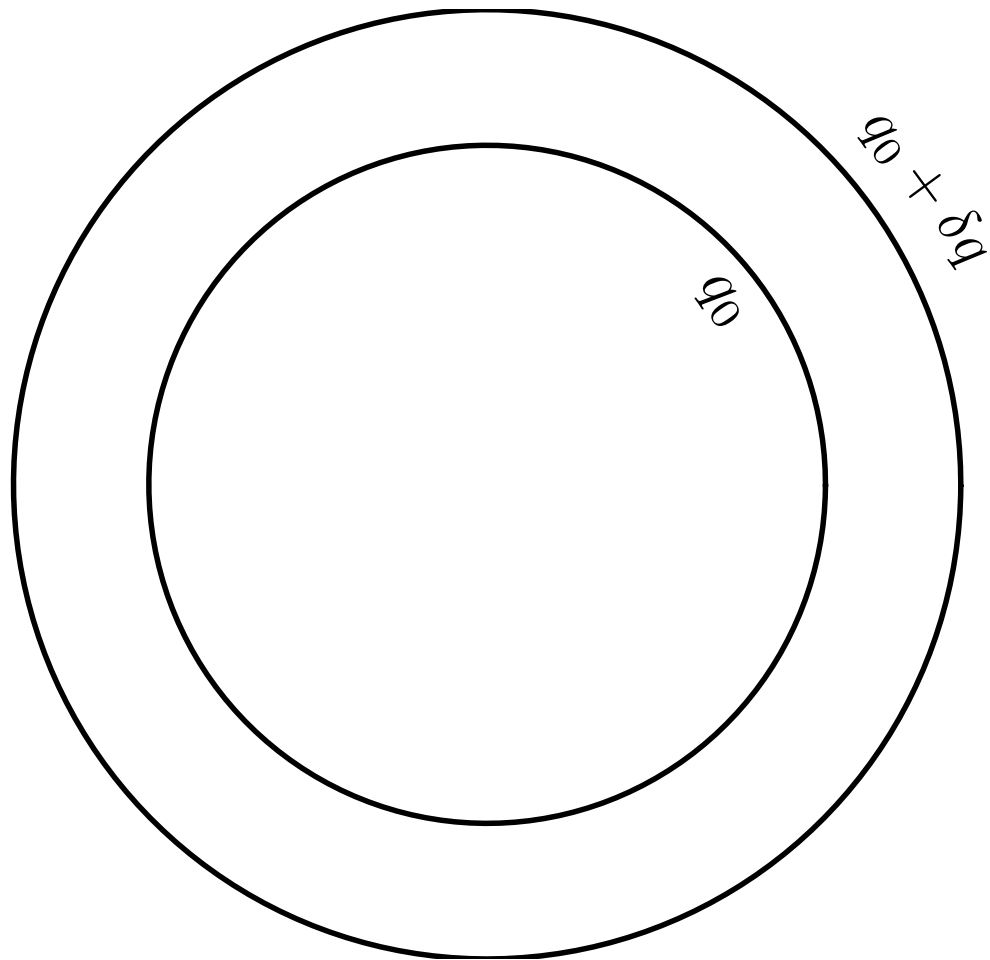
tracer variance forward cascade



2D turbulence

tracer variance forward cascade

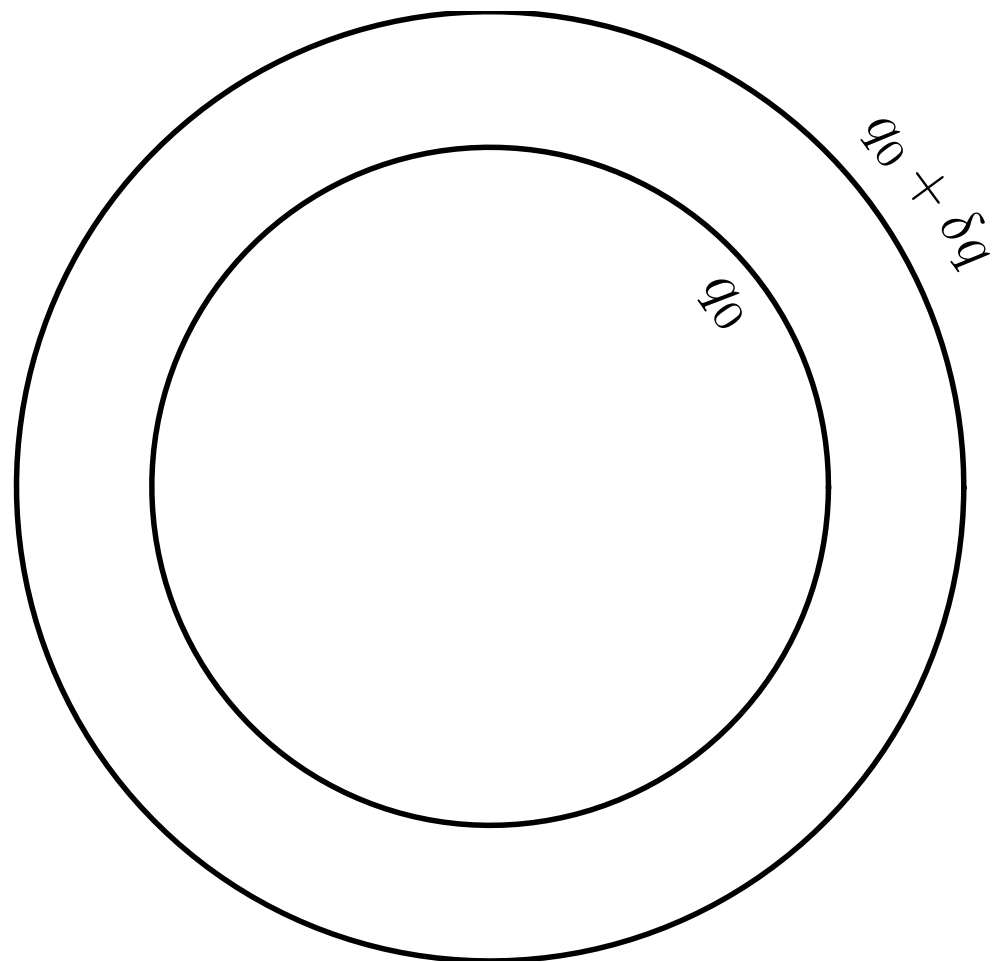
Initial condition



2D turbulence

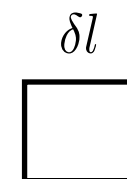
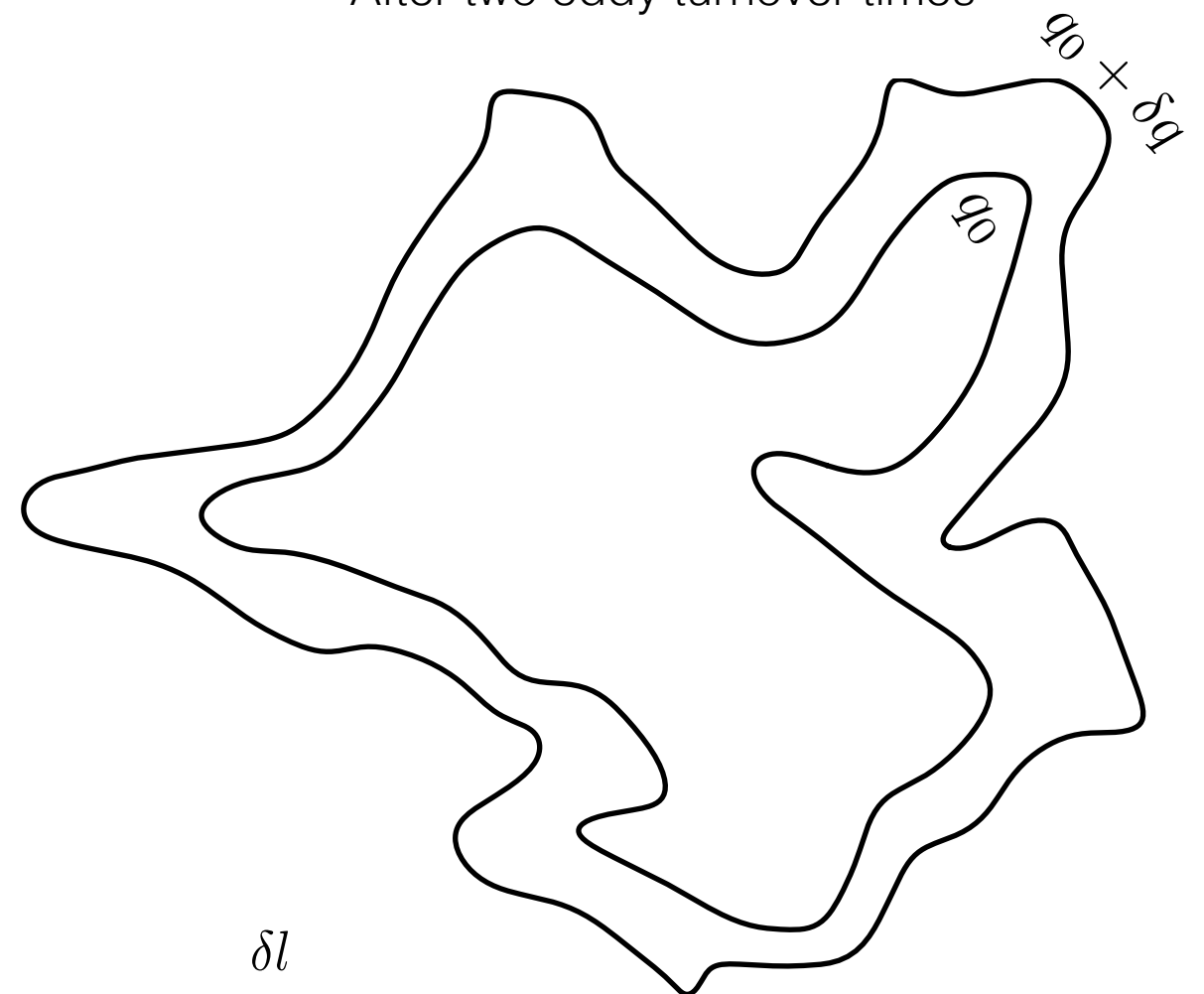
tracer variance forward cascade

Initial condition



Conservation of area

After two eddy turnover times



$$\delta n = \delta q / |\nabla q|$$

(This comes from a 2D stirring simulation with a velocity field represented by “renovating waves”)

2D turbulence

enstrophy forward cascade

The evolution of the active tracer q is more complicated because two quadratic quantities must be conserved in the inviscid limit:

(Kinetic) Energy $E = \frac{1}{2} \iint |\nabla \psi|^2 dx dy = \sum_{|k|} E(|k|)$

(Potential) Enstrophy $Q = \frac{1}{2} \iint q^2 dx dy = \sum_{|k|} |k|^2 E(|k|) + \dots$

$$q = \nabla^2 \psi + h$$

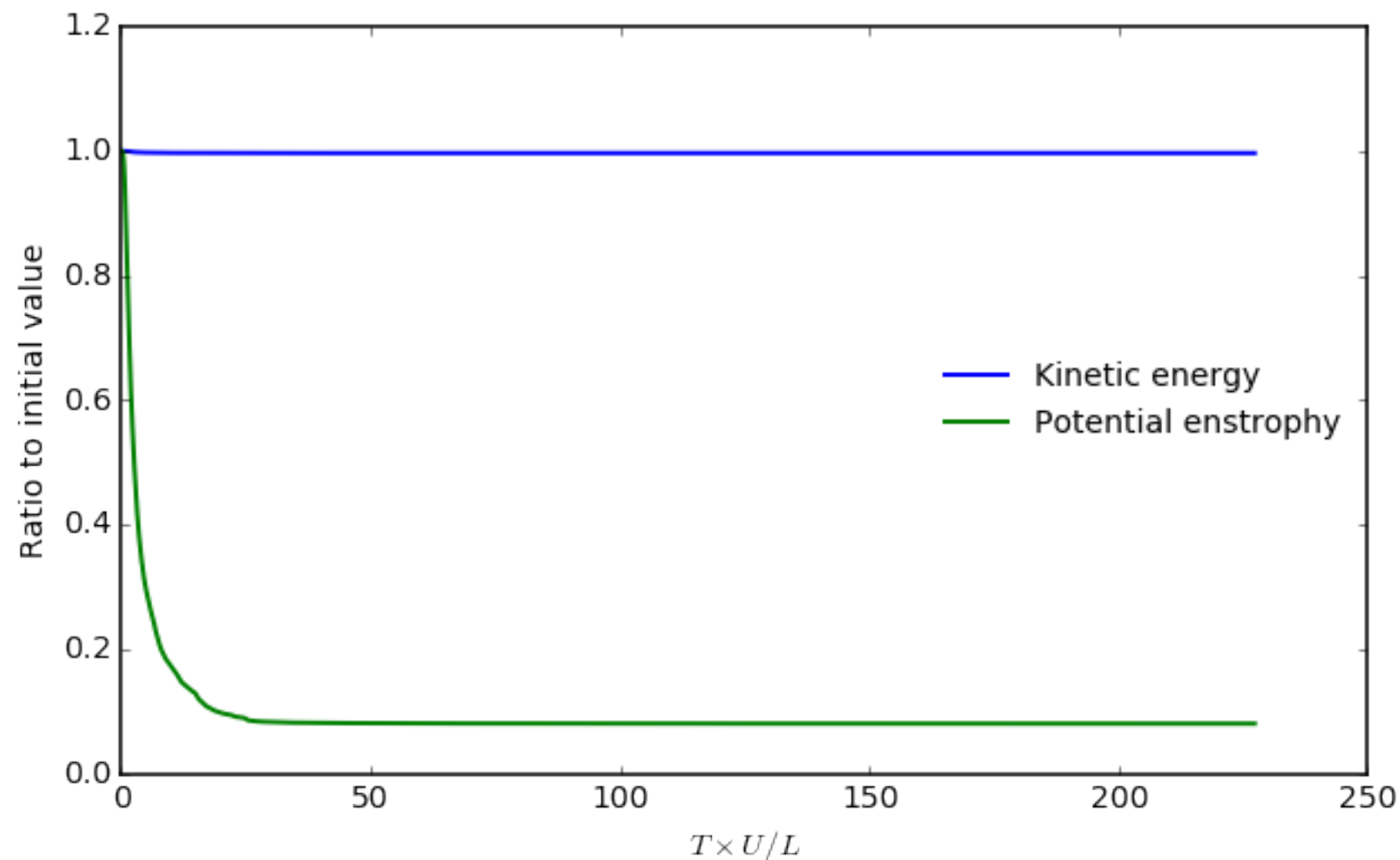
$$|k| = \kappa = (k^2 + l^2)^{1/2}$$

Small-scale dissipation damps enstrophy much more efficiently



2D turbulence: enstrophy decay

Time series in an initial value problem with $h=0$



Energy has only a 0.3 % decay, whereas enstrophy decays by 92%

A minimum enstrophy
principle

What is the flow that minimizes potential enstrophy Q given the energy E ?

“A simple exercise in calculus of variations” (BH76):

$$\begin{aligned}\delta Q + \mu \delta E &= \iint [q \nabla^2 \delta \psi + \mu (\nabla \psi \cdot \nabla \delta \psi)] \, dx \, dy \\ &= \iint \nabla^2 (\nabla^2 \psi + h - \mu \psi) \delta \psi \, dx \, dy = 0\end{aligned}$$

Assuming periodic BCs:

$$(\mu - \nabla^2) \psi = h$$

$$E = \frac{1}{2} \iint |\nabla \psi|^2 \, dx \, dy \qquad Q = \frac{1}{2} \iint q^2 \, dx \, dy$$

What is the flow that minimizes potential enstrophy Q given the energy E ?

Physical space

$$(\mu - \nabla^2) \psi = h$$

Fourier space

$$\hat{\psi}_0 = \frac{\hat{h}}{k^2 + l^2 + \mu}$$

The Lagrangian multiplier defines a length scale $L_0 = \mu^{-1/2}$

$|k|L_0 \ll 1$: $\psi_0 \approx \mu^{-1}h$ Isobaths are streamlines of the coarse-grained flow

$|k|L_0 \gg 1$: $q = \nabla^2\psi_0 + h \approx 0$ The fine-grained PV is homogenized

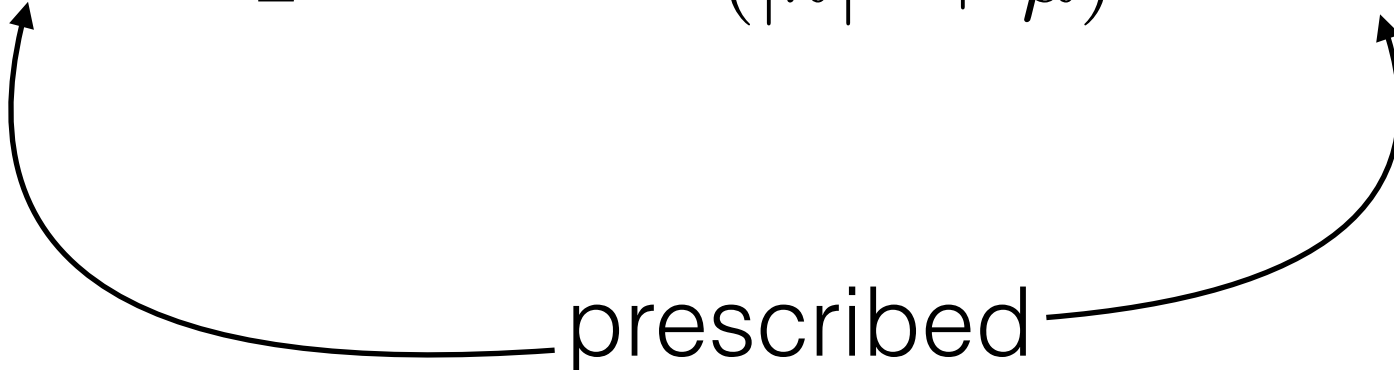
$$|k| = \kappa = (k^2 + l^2)^{1/2}$$

What is the flow that minimizes potential enstrophy Q given the energy E ?

$$\hat{\psi}_0 = \frac{\hat{h}}{k^2 + l^2 + \mu}$$

μ is determined using the constraint E

$$E = \frac{\overline{U^2}}{2} = \frac{1}{2} \sum \frac{|k|^2}{(|k|^2 + \mu)^2} |\hat{h}|^2(|k|)$$



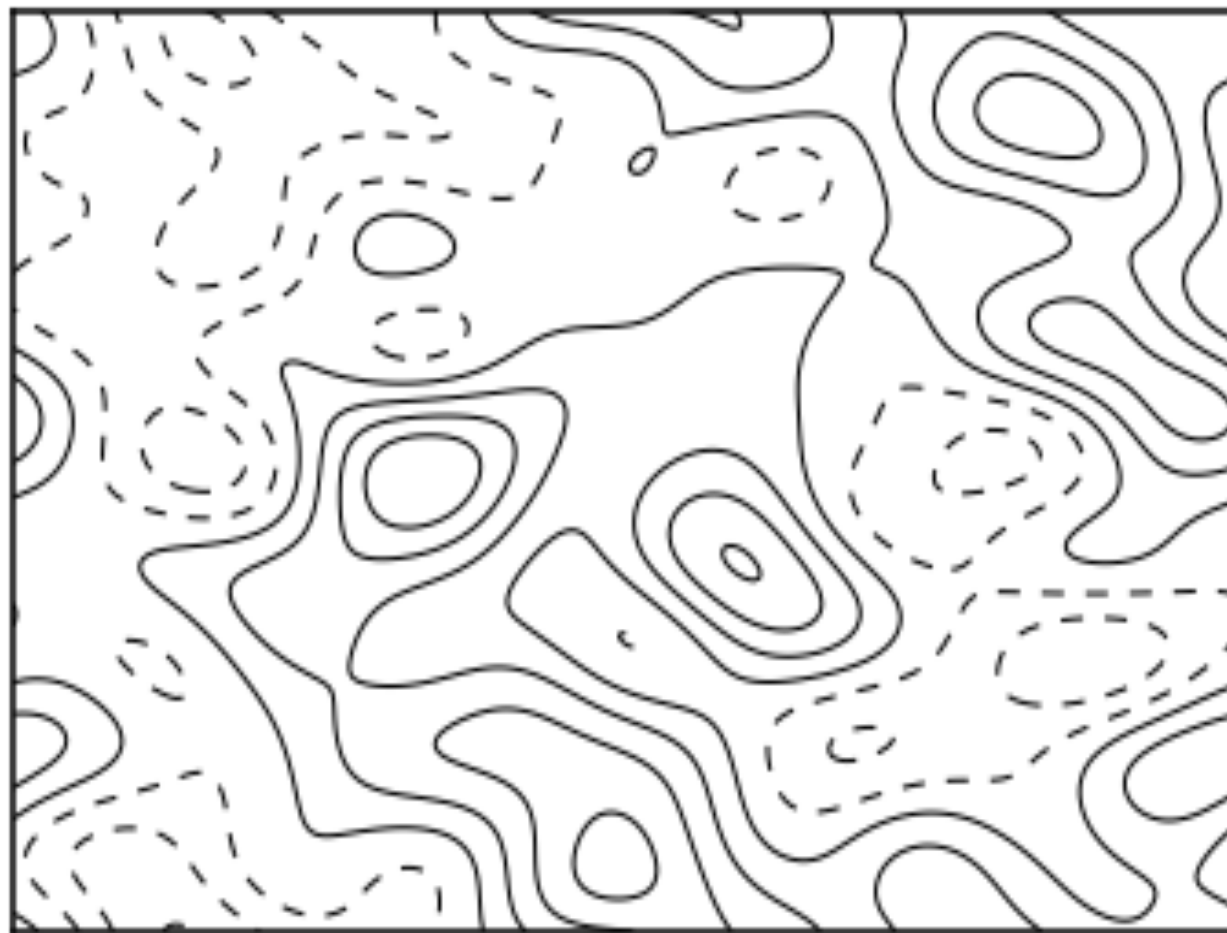
$|k| = \kappa = (k^2 + l^2)^{1/2}$

prescribed

An example

“The preferred topography” (BH76)

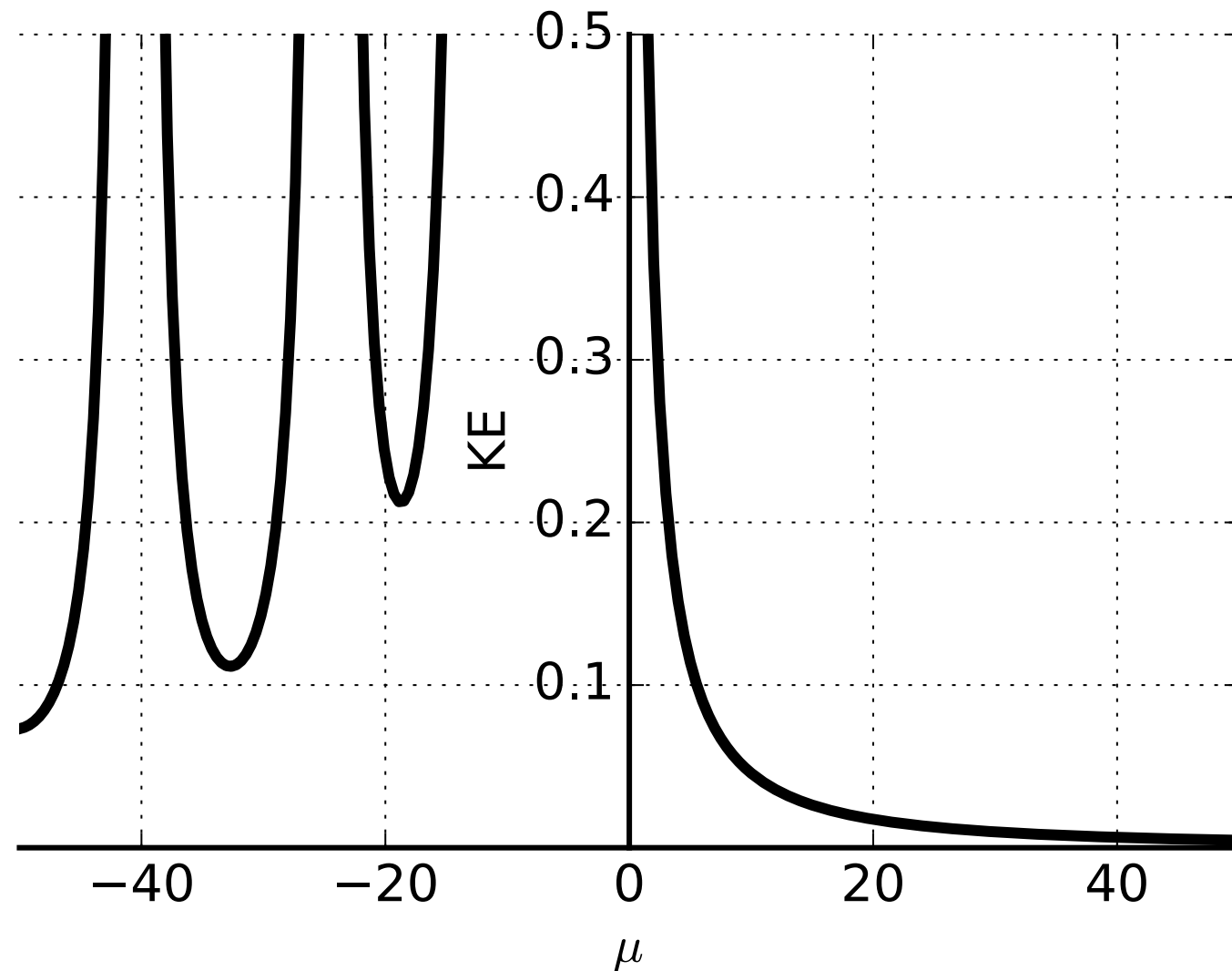
$$\hat{h}(|k|) \propto |k|^{-2} \quad |k|_{min} = 1, |k|_{max} = 12$$



$$|k| = \kappa = (k^2 + l^2)^{1/2}$$

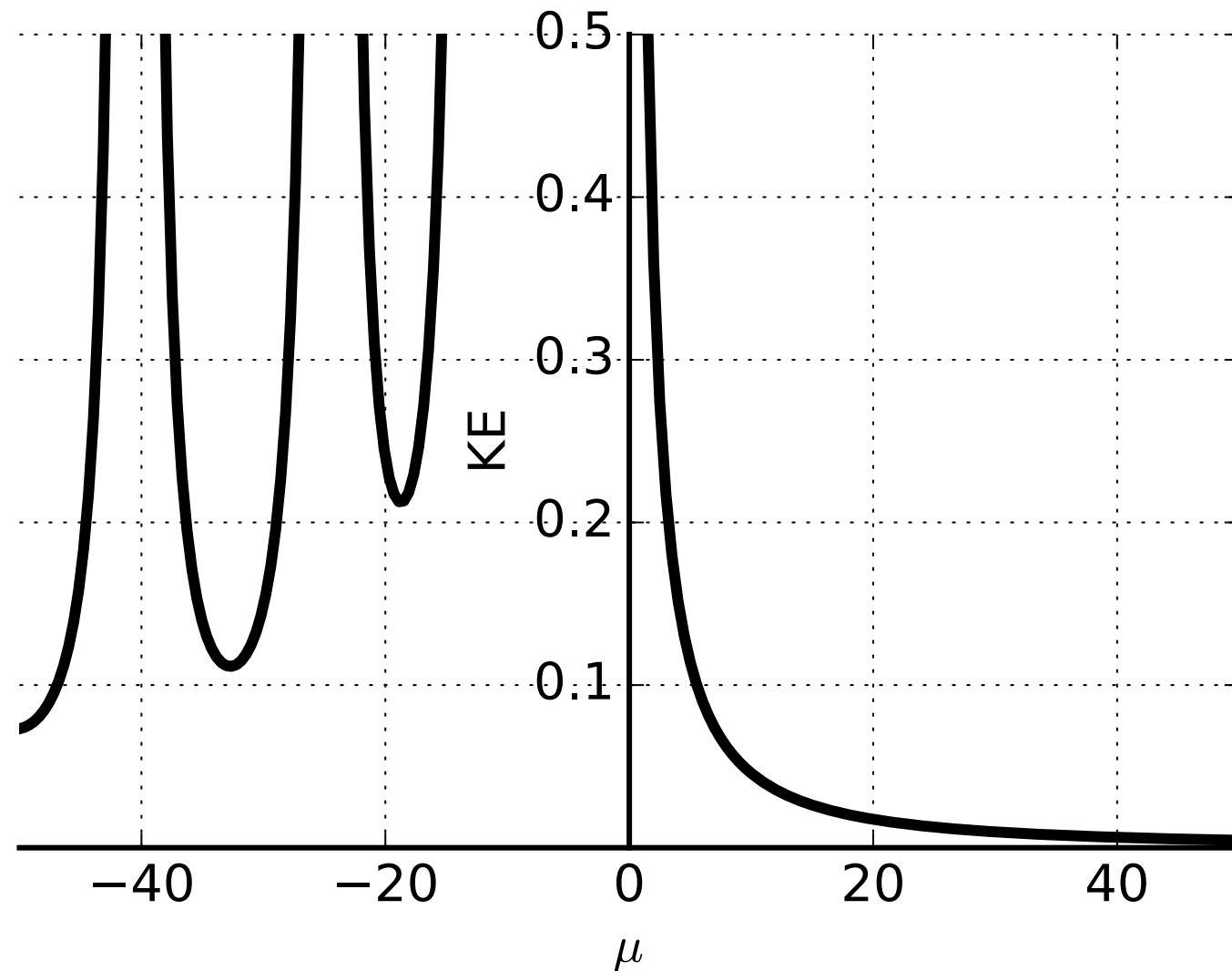
$$|k|_{min} = 1, |k|_{max} = 12$$

A given energy level has multiple solutions
but the positive one minimizes enstrophy



“It is readily demonstrated that for positive μ this is in fact a minimum” (BH76)

A given energy level has multiple solutions
but the positive one minimizes enstrophy



The second order terms in the variational problem

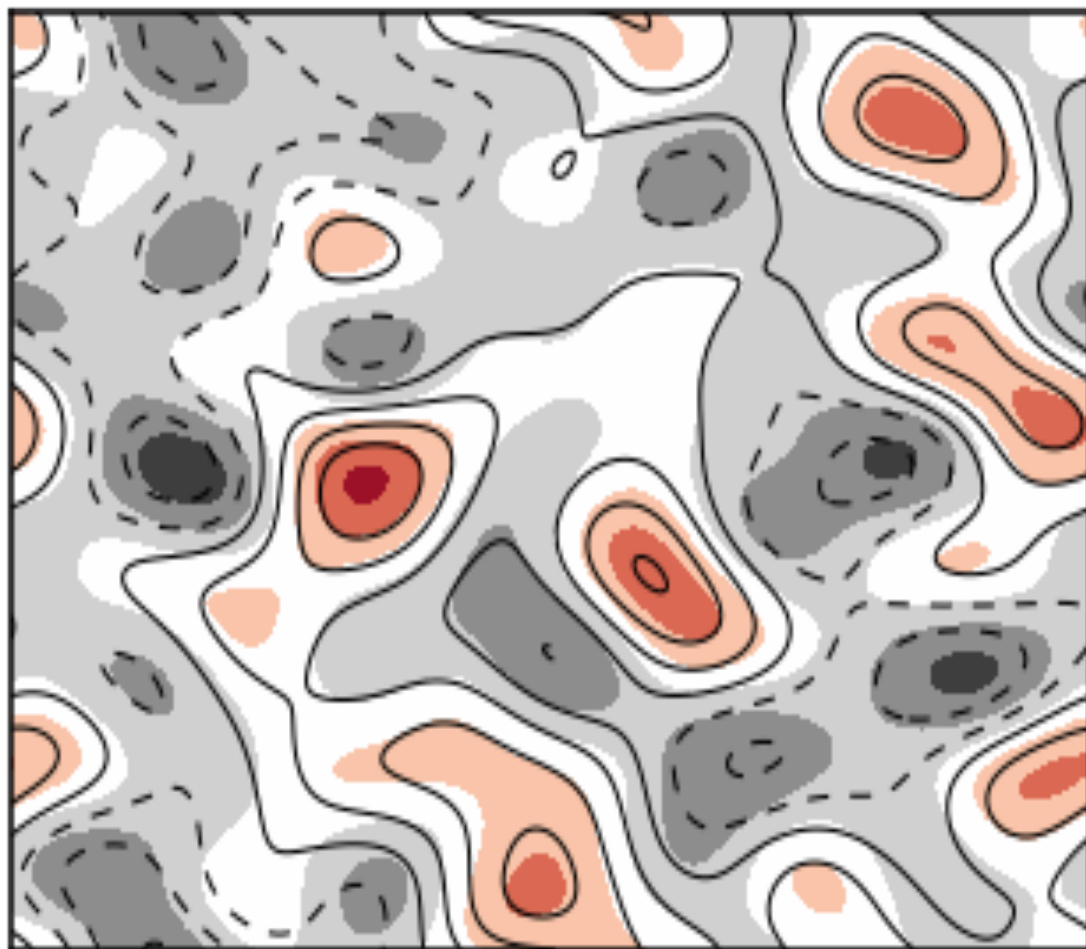
$$\delta^2 Q + \mu \delta^2 E = \frac{1}{2} \iint (\nabla^2 \delta \psi)^2 dx dy + \frac{\mu}{2} \iint |\nabla \delta \psi|^2 dx dy$$

An example of minimum enstrophy solution

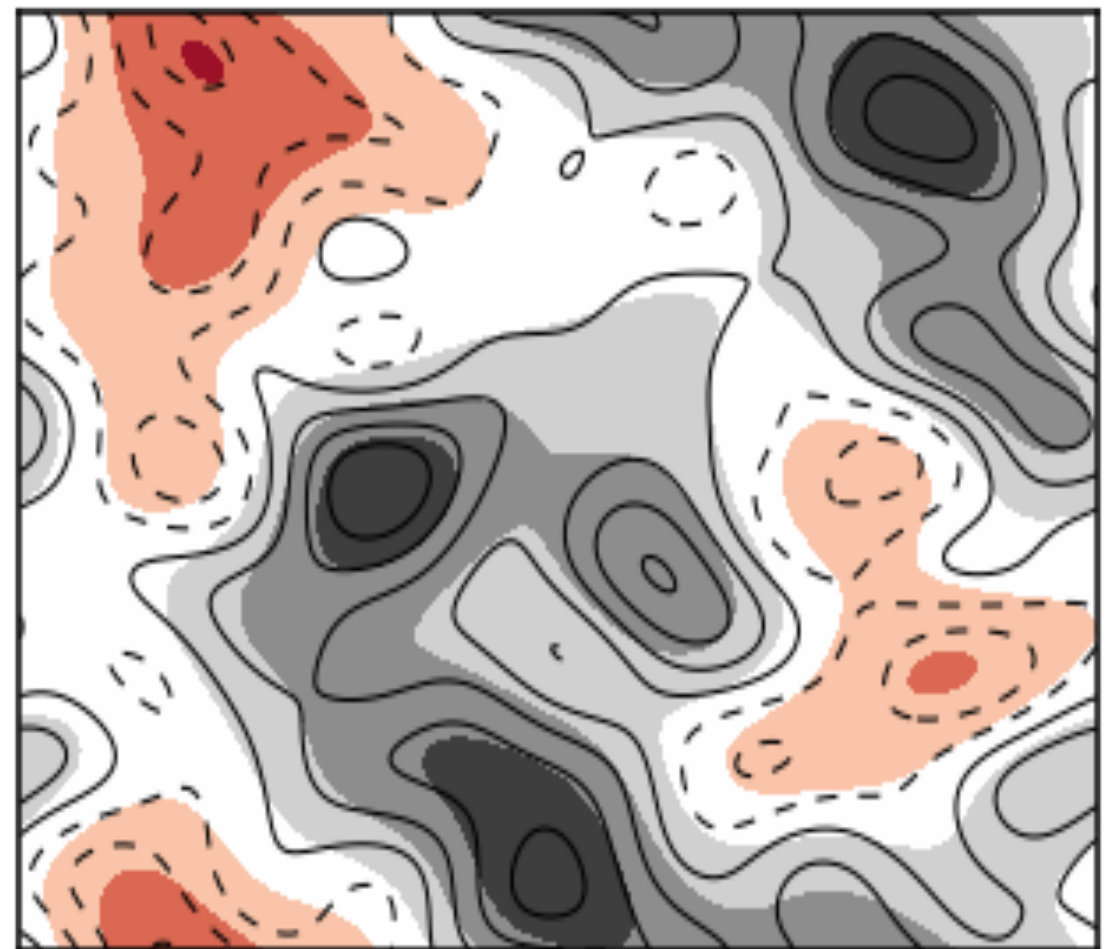
$$E = 0.05, \quad \mu = 9.22072072$$

contours: h

Colors: $PV \nabla^2 \psi_0 + h$



Colors: Streamfunction ψ_0

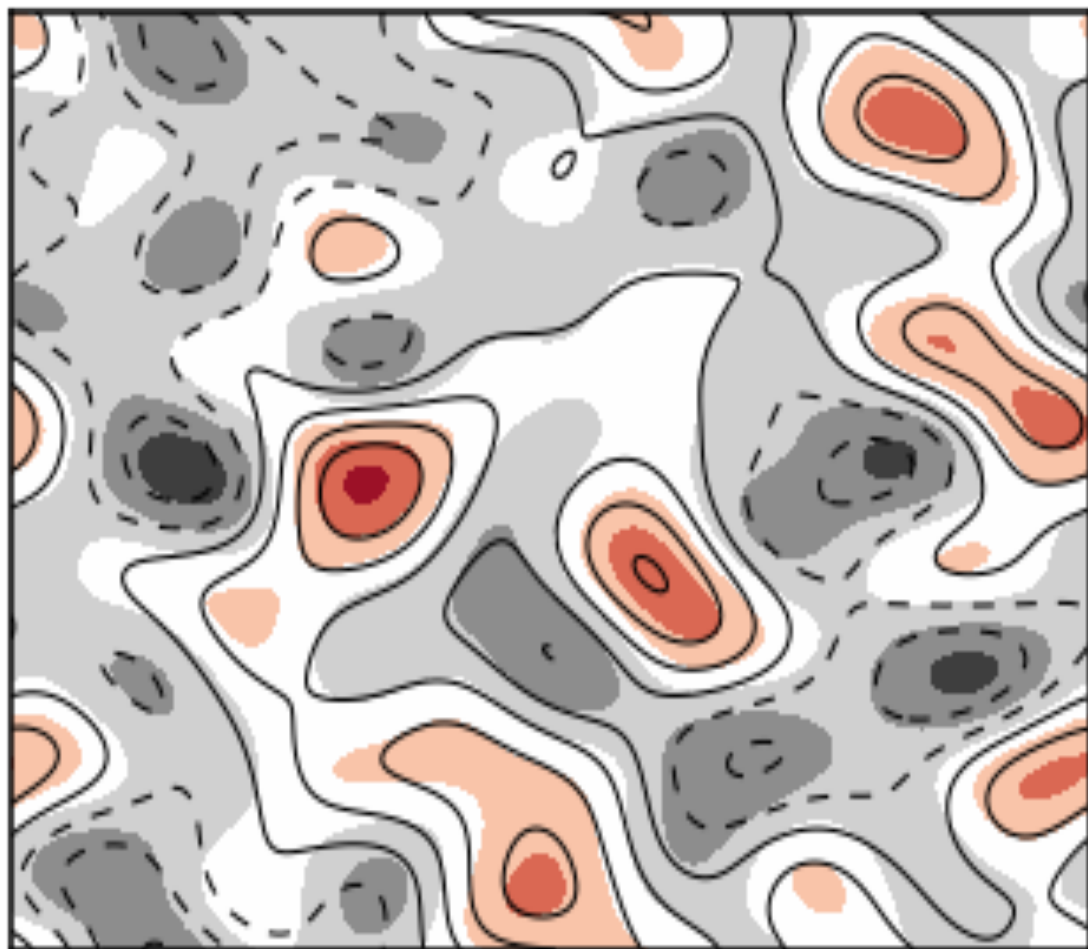


An example of minimum enstrophy solution

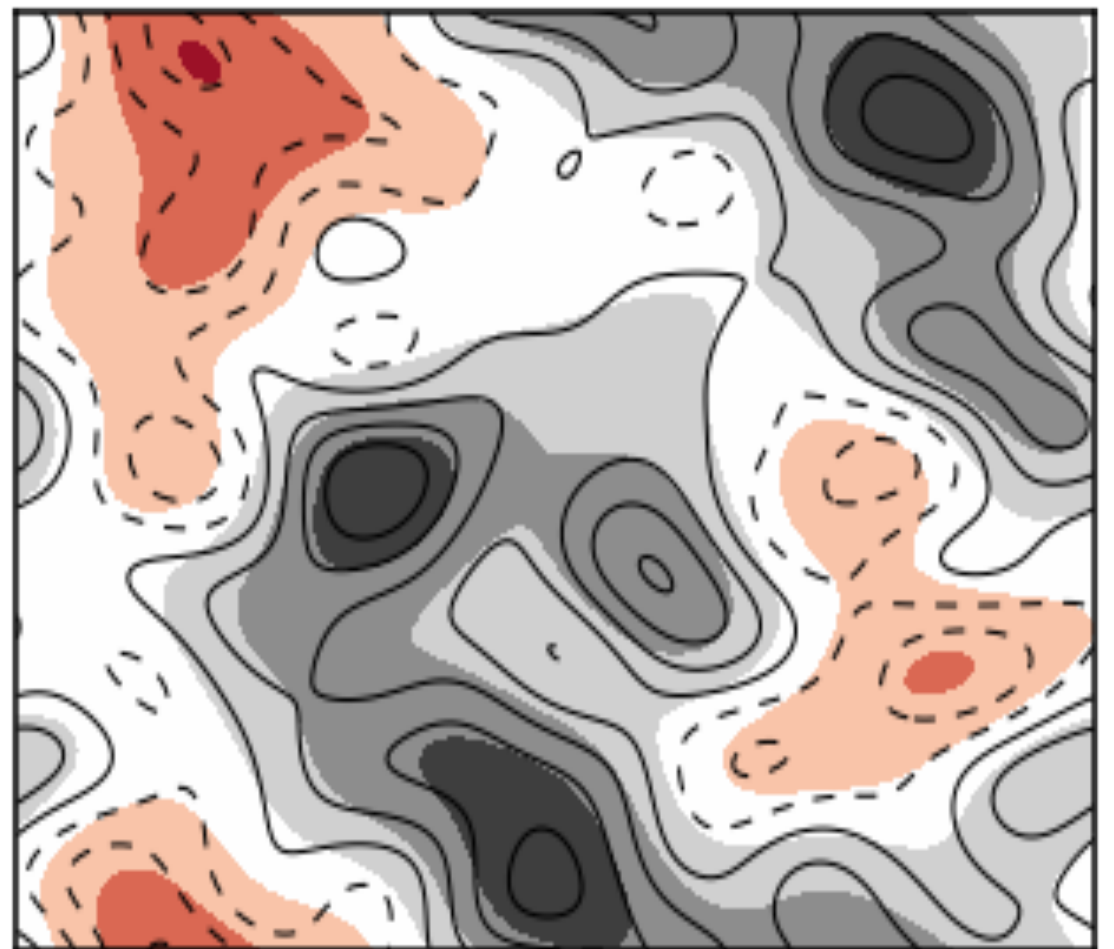
$$E = 0.05, \quad \mu = 9.22072072$$

contours: h

Colors: $PV \nabla^2 \psi_0 + h$



Colors: Streamfunction ψ_0



Discuss: Why BH76 did not do this!?

I will not discuss.

The role of saddle points

(There a lots of typos in equations 24 through 30 of the paper.)

The role of viscosity

The role of viscosity

$$\frac{Dq}{Dt} = \nu \nabla^2 (\nabla^2 \psi)$$

$$q = \nabla^2 \psi + h$$

$$\begin{aligned} \frac{\partial Q}{\partial t} &= -\nu \iint \nabla q \cdot \nabla (q - h) dx dy \\ &= -\nu \iint |\nabla q|^2 dx dy + \nu \iint \nabla q \cdot \nabla h dx dy \end{aligned}$$

↑

Positive or negative?

If, $\psi = \psi_0$ then enstrophy must increase since ψ_0 is the minimum enstrophy solution.

Closed-basin solution

The circulation on the boundary must be prescribed

$$\text{If } \psi = 0 \text{ on } \Gamma \text{ then } C = \oint_{\Gamma} \partial_n \psi \, ds = \text{constant}$$

The minimum enstrophy problem has two constraints

$$\delta \frac{1}{2} \iint (\nabla^2 \psi + h)^2 \, dx \, dy + \mu \delta \frac{1}{2} \iint |\nabla \psi|^2 \, dx \, dy + \lambda \oint \frac{\partial}{\partial n} \delta \psi \, ds = 0$$

Use integration by parts, e.g.,

$$\iint \nabla^2 \psi \nabla^2 \delta \psi \, dx \, dy = \underbrace{\iint \nabla \cdot [\nabla^2 \psi \nabla \delta \psi] \, dx \, dy}_{= \oint \nabla^2 \psi \frac{\partial}{\partial n} \delta \psi \, ds} - \underbrace{\iint \nabla \nabla^2 \psi \cdot \nabla \delta \psi \, dx \, dy}_{= + \iint (\nabla^2 \nabla^2 \psi) \delta \psi \, dx \, dy}$$

The circulation on the boundary must be prescribed

$$\text{If } \psi = 0 \text{ on } \Gamma \text{ then } C = \oint_{\Gamma} \partial_n \psi \, ds = \text{constant}$$

The minimum enstrophy problem has two constraints

$$\nabla^2(\nabla^2 \psi + h - \mu \psi) = 0 \quad \text{within } S$$

$$\nabla^2 \psi + h + \lambda = 0 \quad \text{on } \Gamma$$

Using $\psi = 0$ on Γ :

$$\nabla^2 \psi - \mu \psi = -(h + \lambda) \quad \text{within } S$$

C determines the new constant λ

A numerical experiment

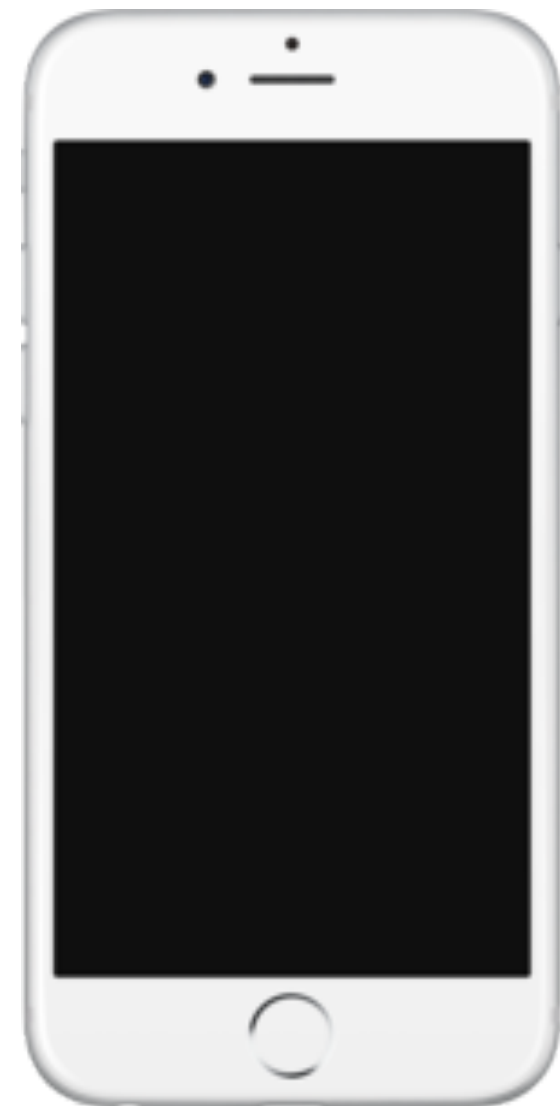
Side remark

CPU 36.4 MHz; 36 MFLOPS



CDC 7600

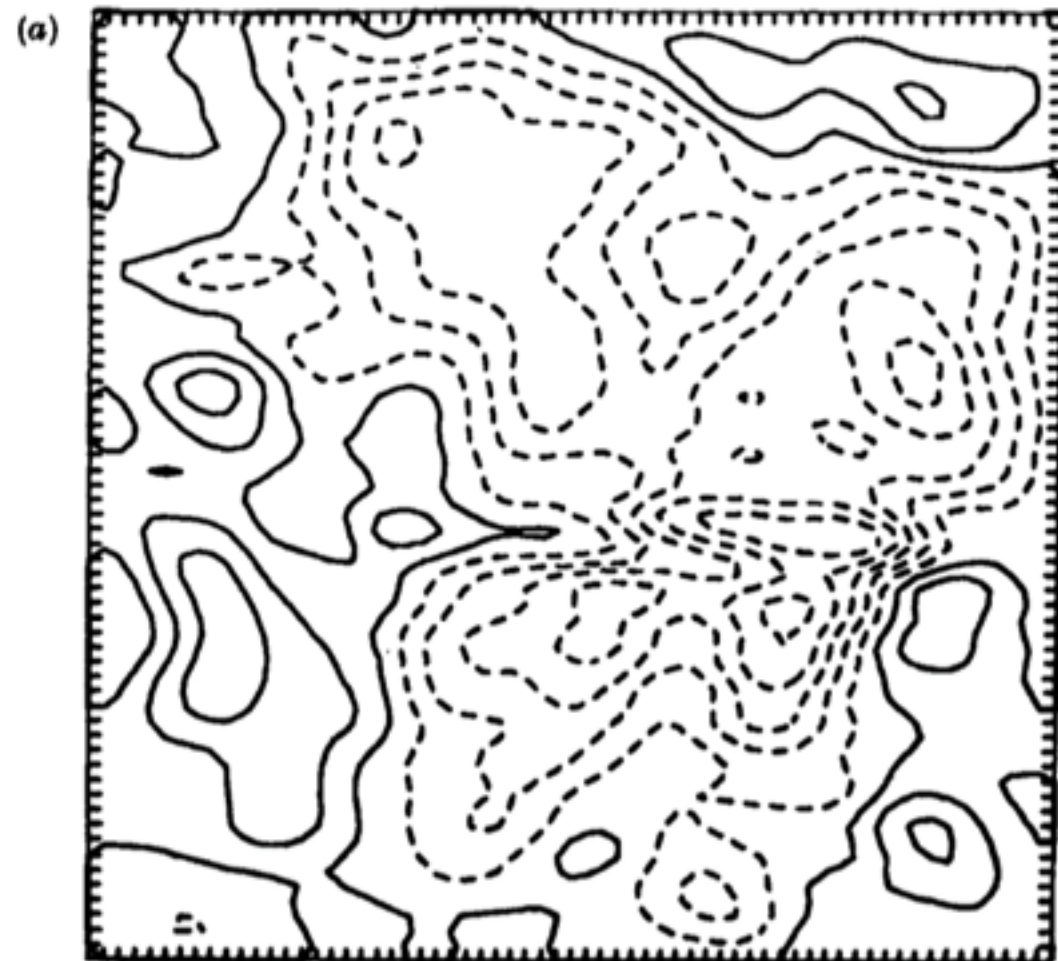
CPU 1.4 GHz; GPU;
100's GFLOPS to TFLOPS



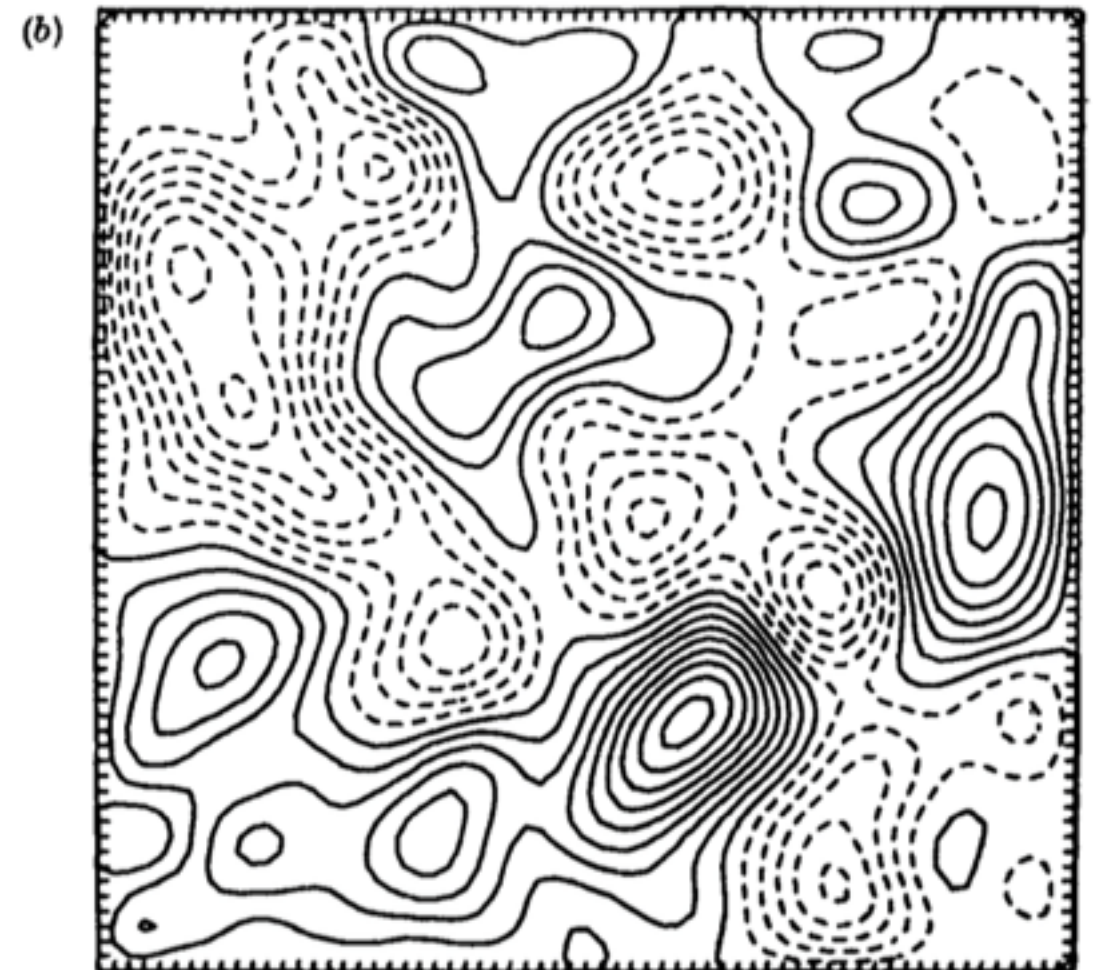
iPhone 6

Source: Wikipedia

Topography

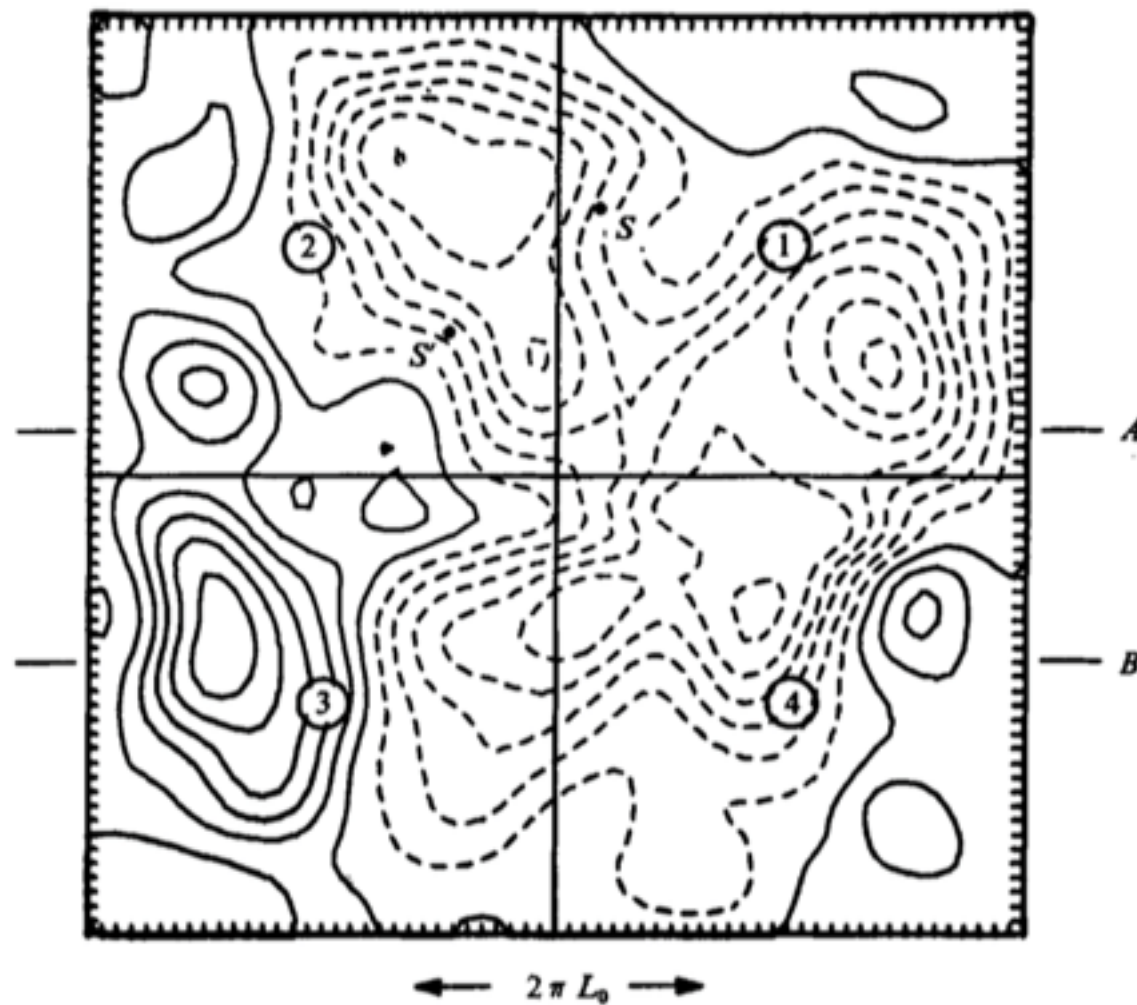


Initial streamfunction

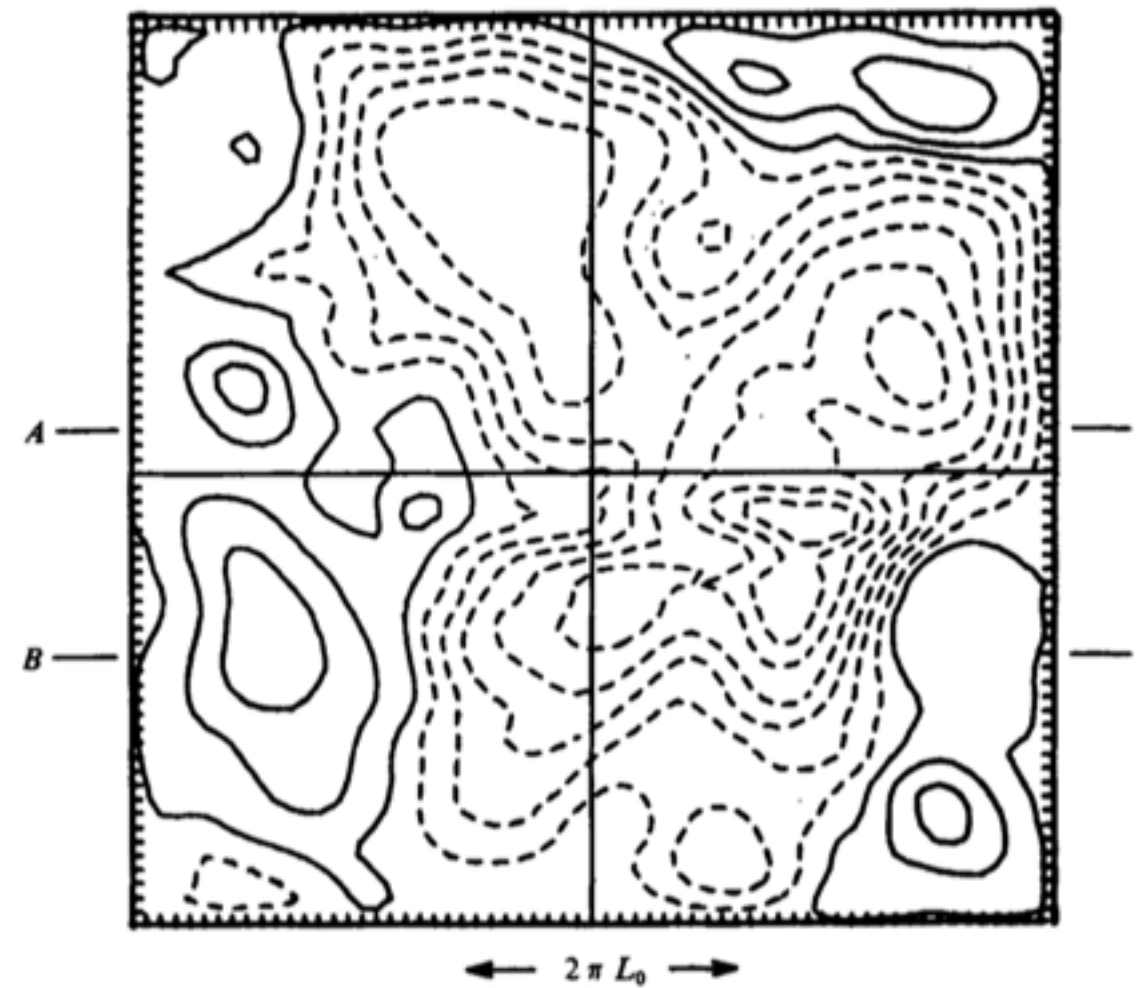


What is the initial energy level and the associated μ ?

Final streamfunction

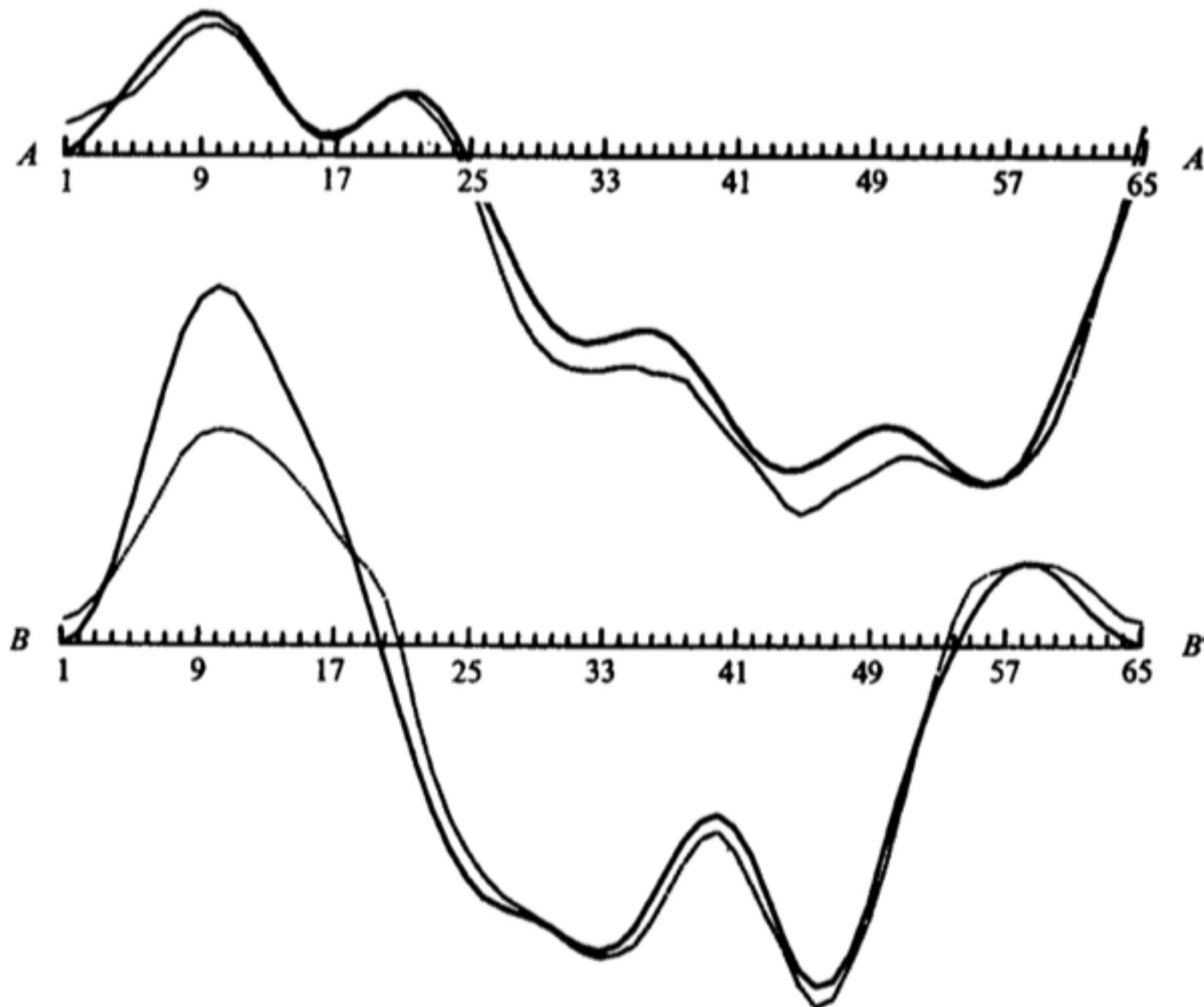


Final PV



The flow quickly becomes quasi-steady

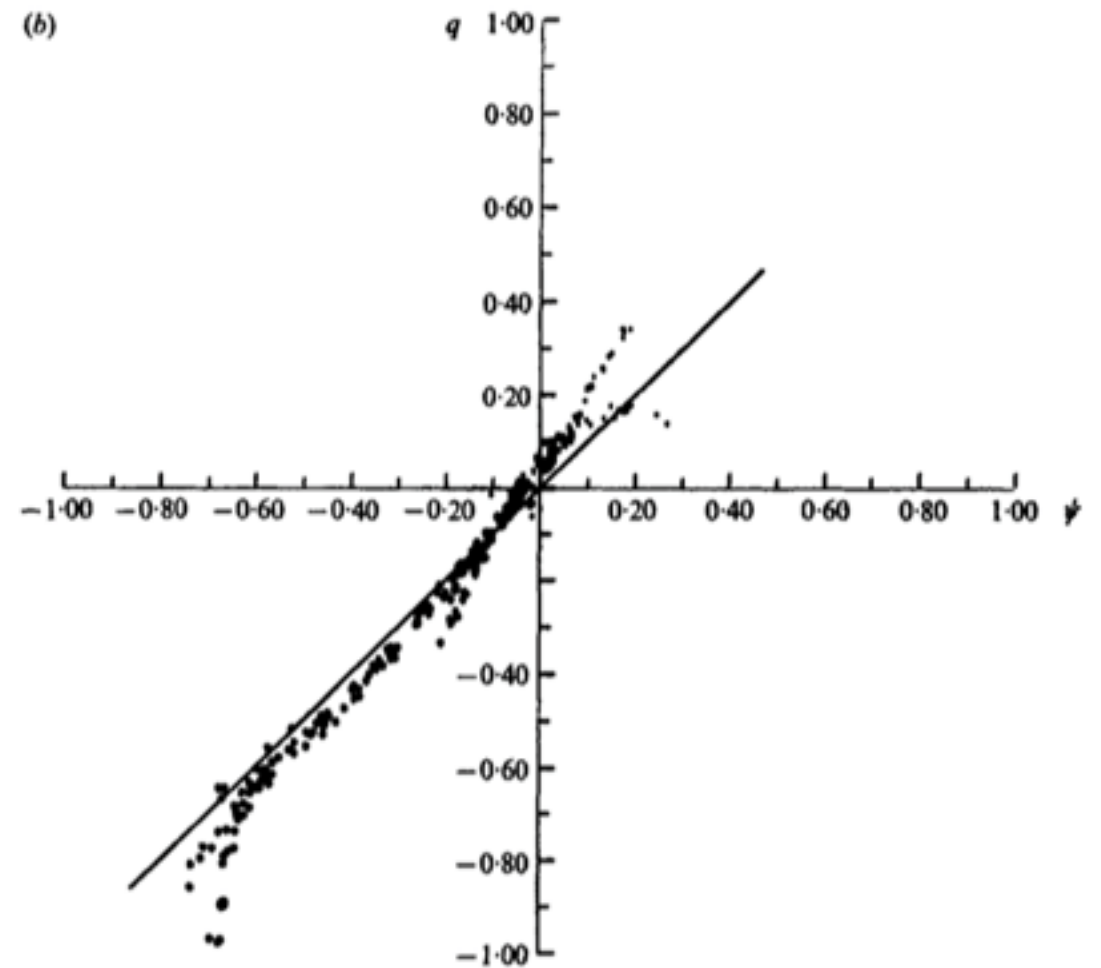
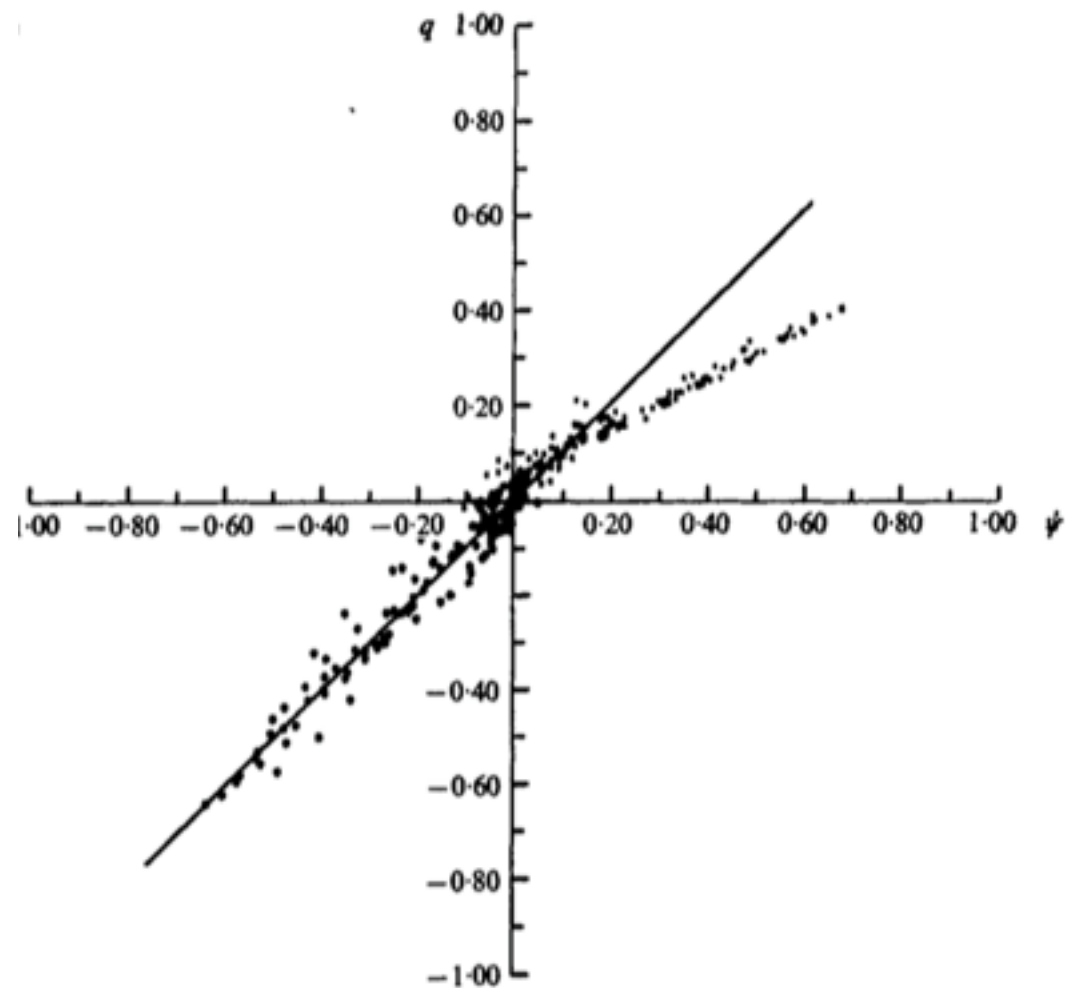
$$J(\psi, q) = 0 \rightarrow q = F(\psi)$$



Streamfunction

PV

The $\psi - q$ relationship

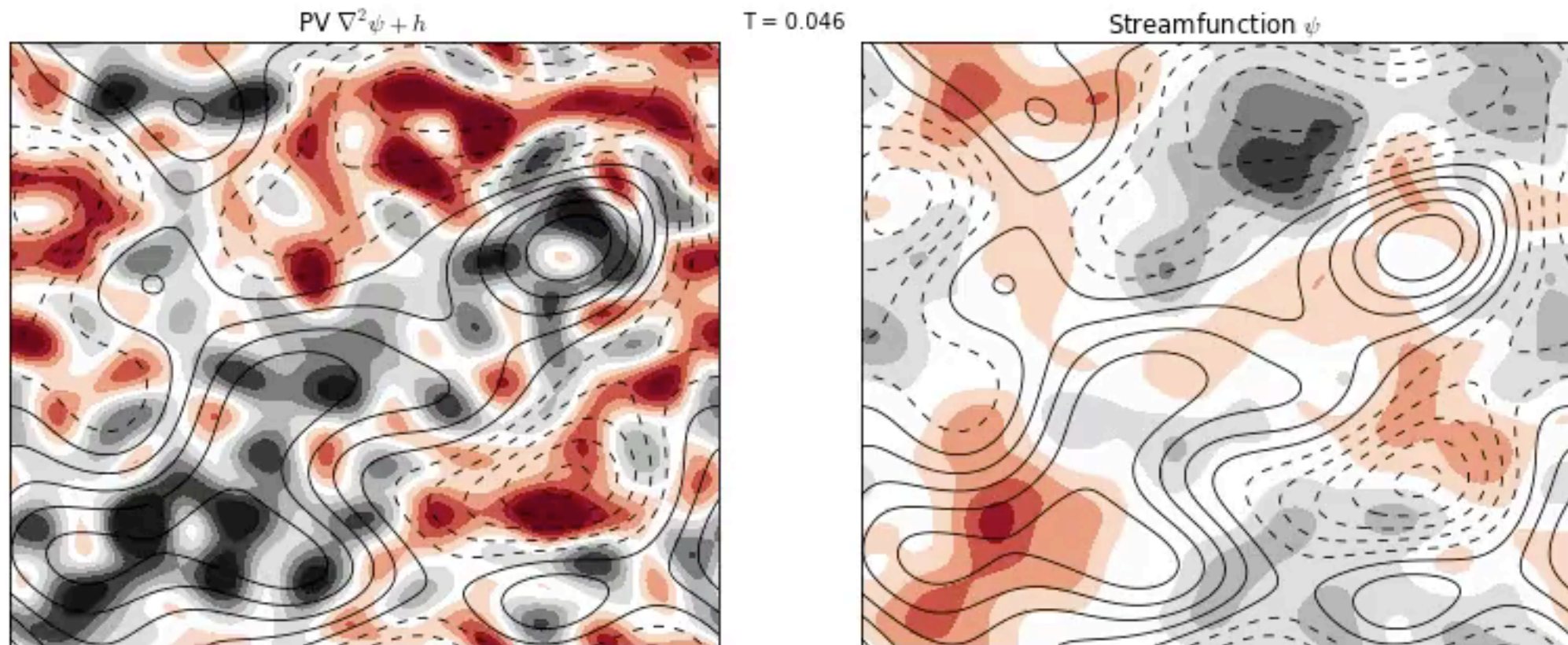


Roughly linear...but there are different “regimes”...

But the energy dropped by
58% owing to low resolution...

$$E = 0.5 \quad |\hat{h}|^2 \propto |k|^{-2} \quad |k|^2 |\psi|^2(t=0) \propto |k| \left[1 + \left(\frac{k}{6} \right)^4 \right]^{-1}$$

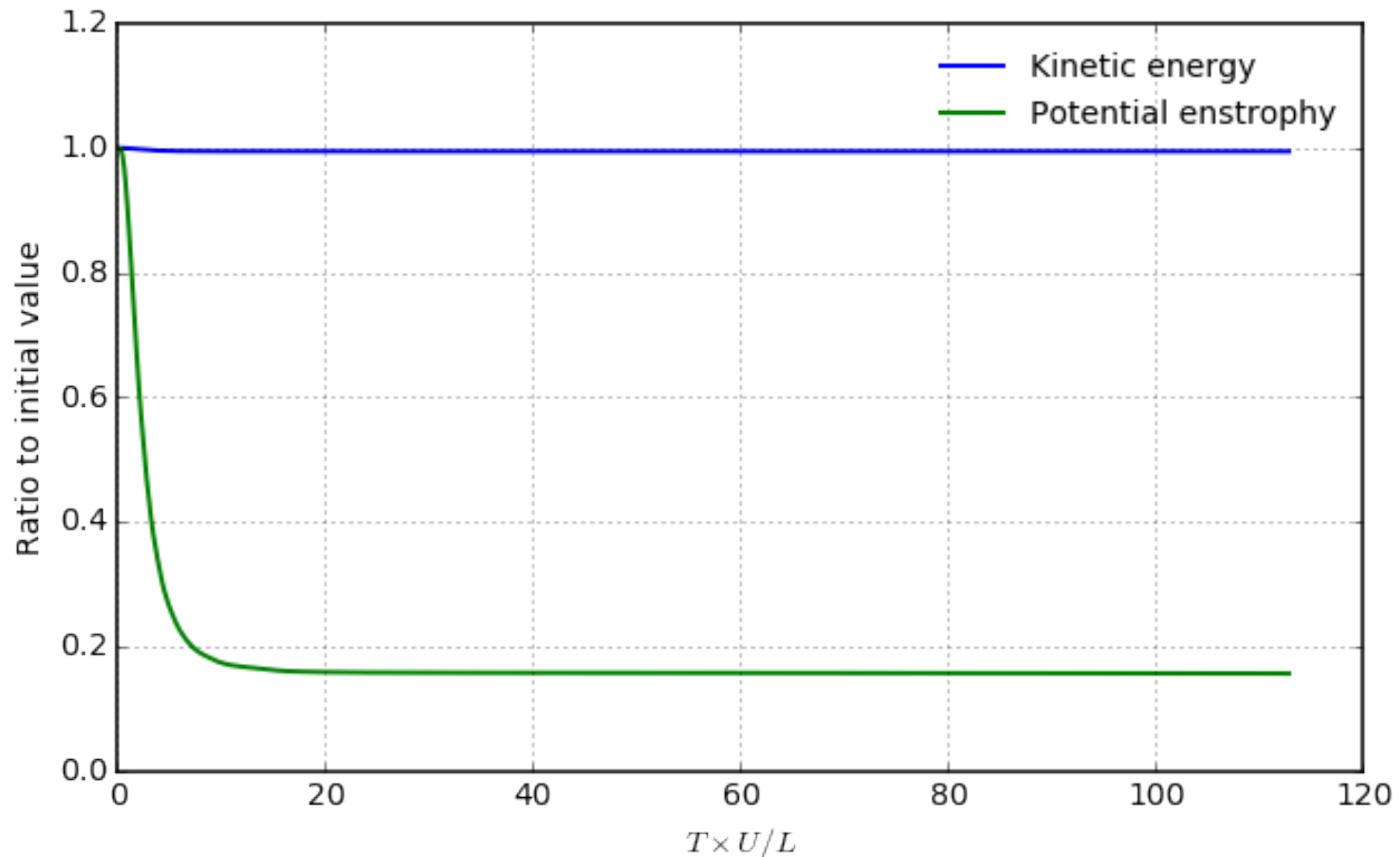
$$\mu = 2.38288288$$



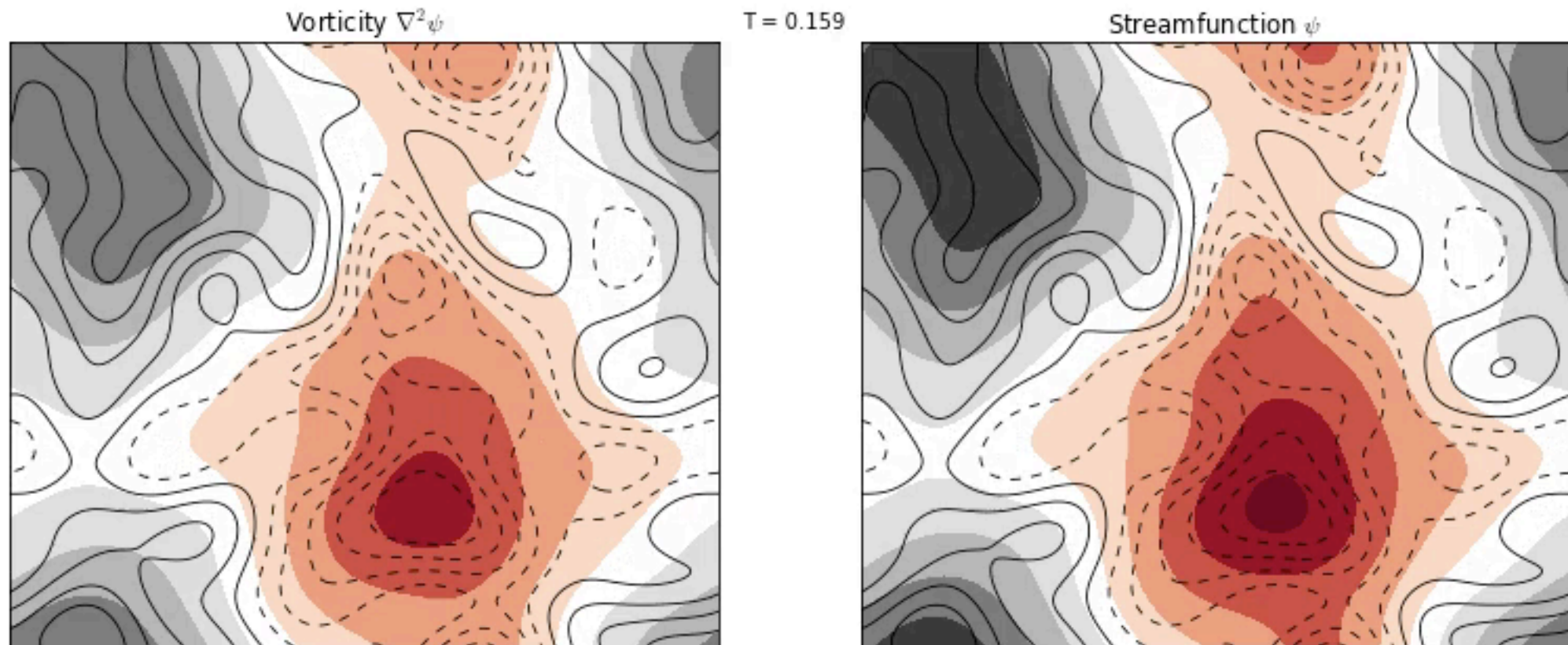
Contours represent topography; colors streamfunction or PV

(doubly periodic calculations)

Energy stays nearly constant,
enstrophy decays significantly...



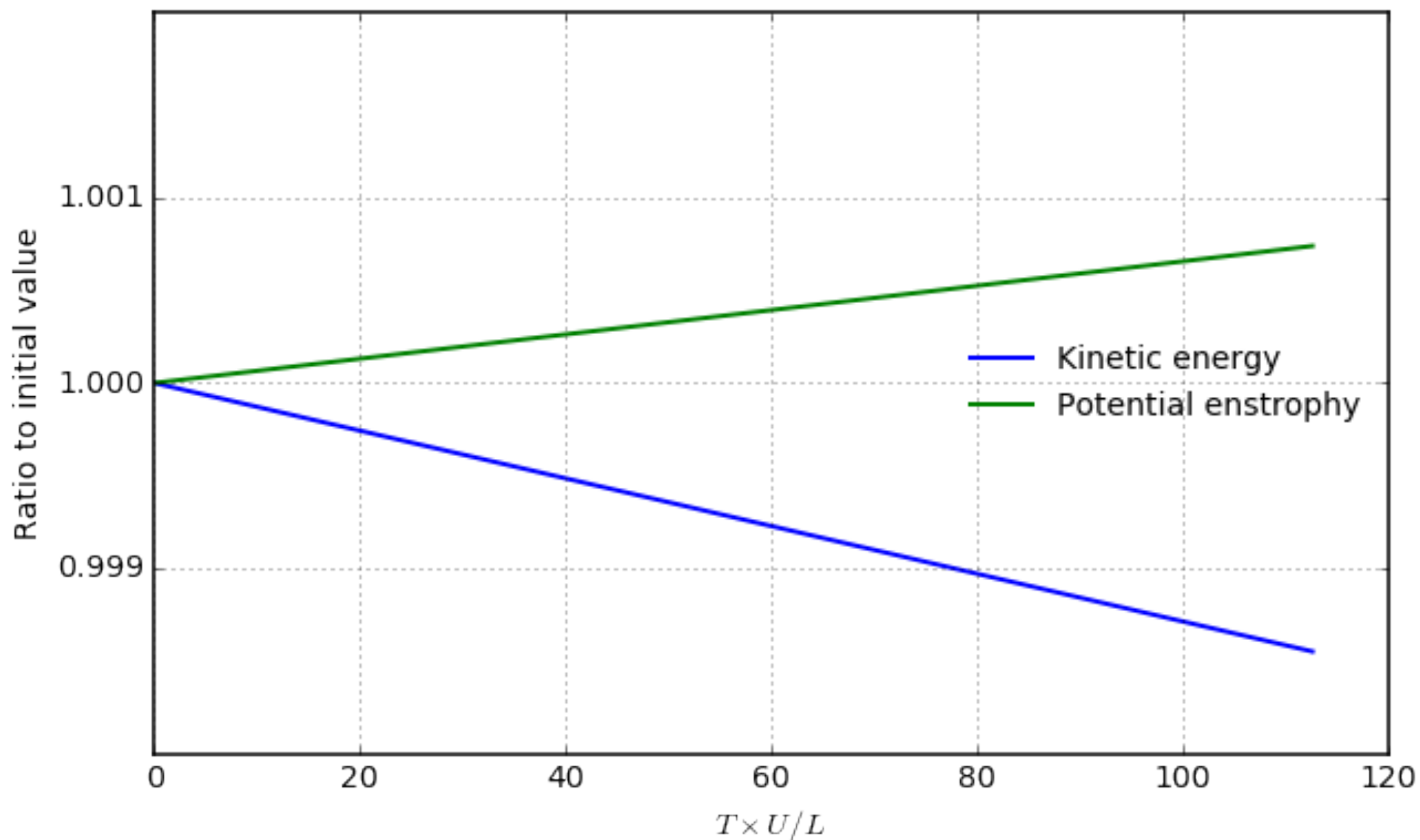
Initializing a simulation with the minimum enstrophy solution



Contours represent topography; colors streamfunction or PV

In practice, with viscosity, energy decays super slowly...

Energy decays, enstrophy increases...



The effects of β

The effects of β

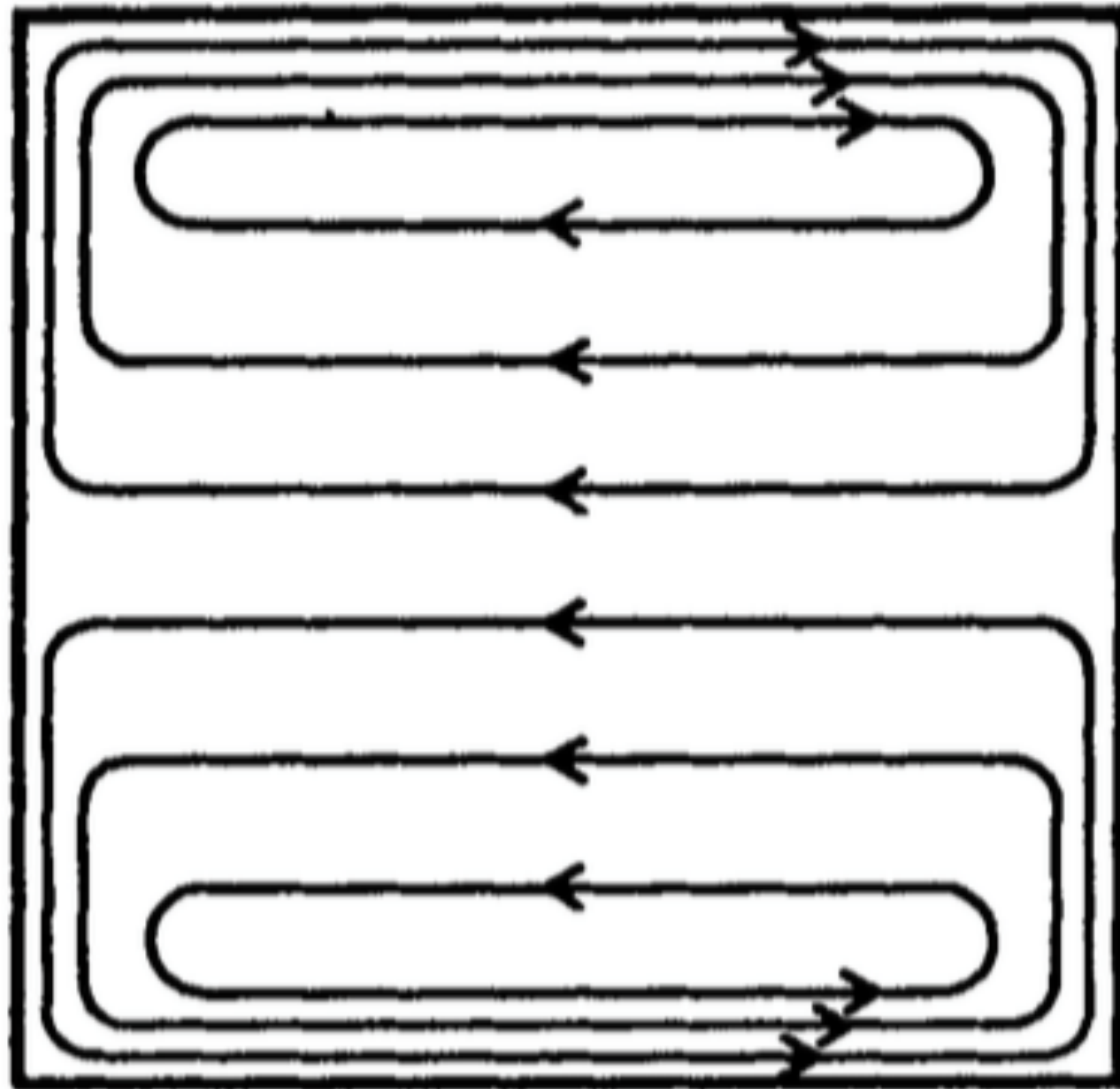
“In so far the the b-slope can be thought of as a Fourier component of zero wavenumber the solution may be instantly obtained”:

$$\psi_0(x, y) = \sum \frac{\hat{h}}{k^2 + l^2 + \mu} + \frac{\beta}{\mu} y$$

Discuss: is this obvious?

The effects of β

Fofonoff mode type of solution

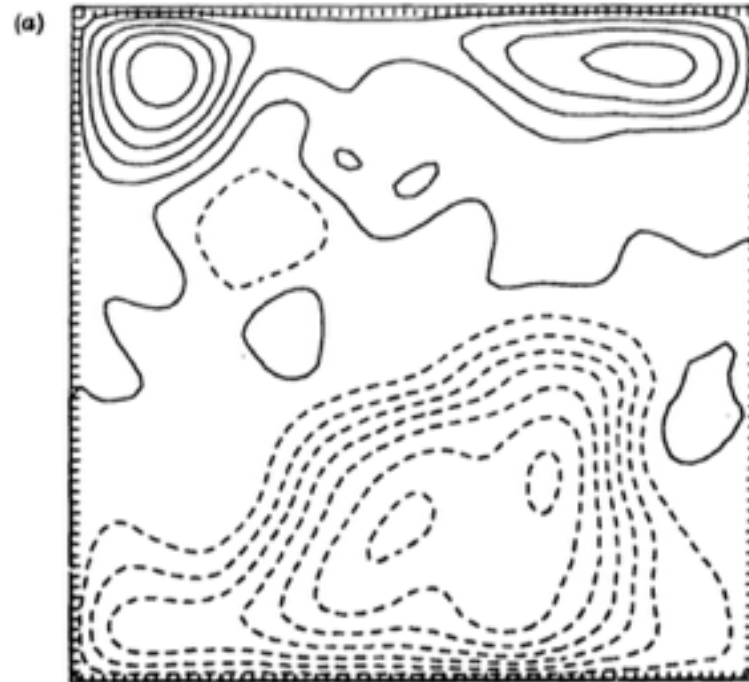


Numerical experiments

Topography



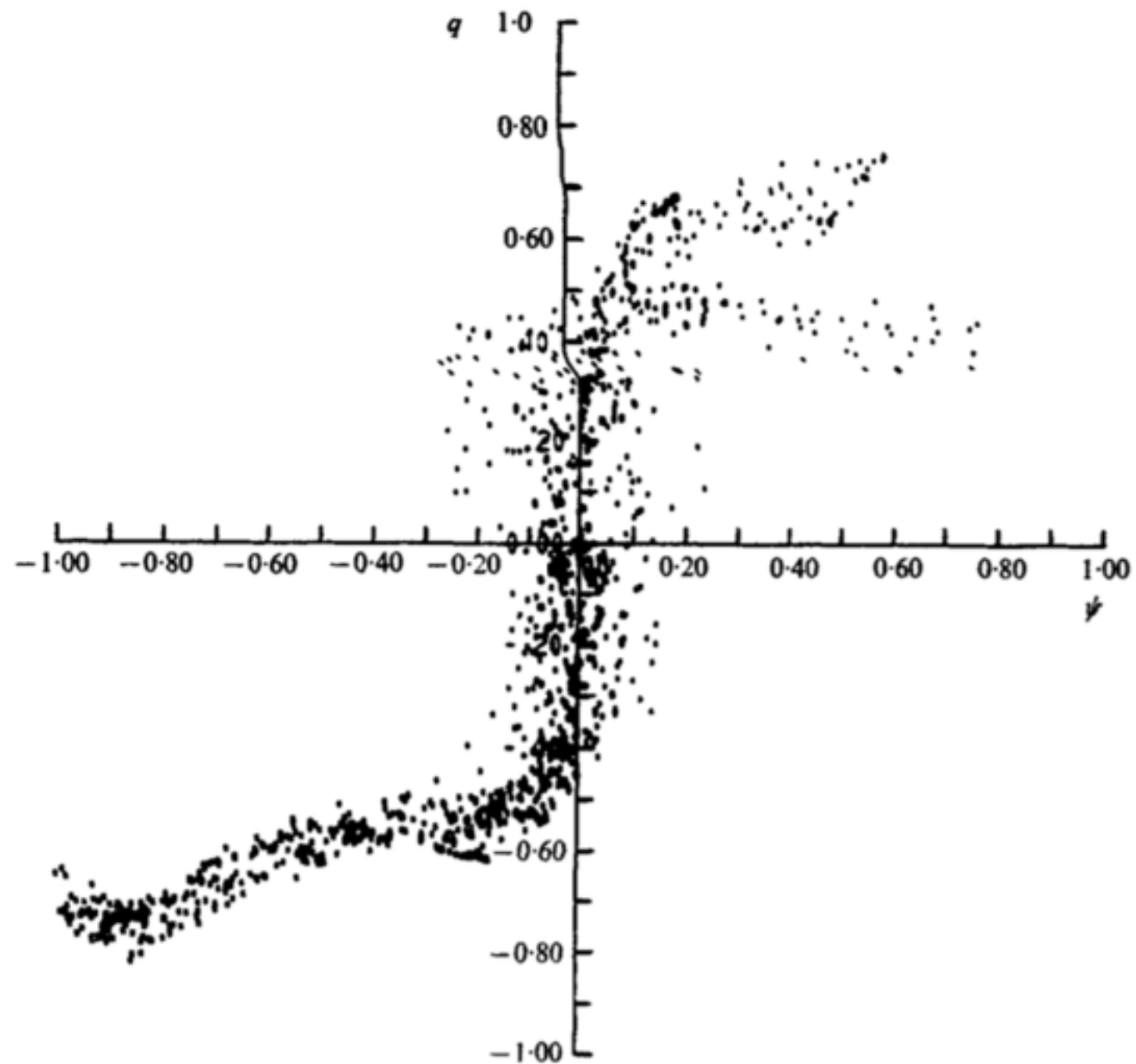
Final
streamfunction



Final PV



The $\psi - q$ relationship



The effects of eddies on the
large scale flow

Eddy fluxes

Reynolds decomposition

$$\psi = \Psi + \psi' \quad q = (\nabla^2 \Psi + \beta y) + \underbrace{(\nabla^2 \psi' + h)}_{q'}$$

.... The eddy PV flux div: $\nabla \cdot \mathbf{F} = \overline{\mathbf{J}(\psi', q')}$

The PV eqn:

$$\frac{Dq'}{Dt} = -\beta \psi'_x$$

Eddy fluxes

$$\frac{Dq'}{Dt} = -\beta\psi'_x$$

$$q' = -\beta\eta$$

η : northward particle displacement

The single-particle diffusivity is

$$\begin{aligned} D &= \frac{d}{dt} \frac{1}{2} \overline{\eta^2} = \beta^{-2} \frac{d}{dt} \frac{1}{2} \overline{q'^2} = \beta^{-1} \overline{v'q'} = -\beta^{-1} \overline{(\nabla^2 \psi' + h)\psi'_x} \\ &\approx -\beta^{-1} \left[\frac{1}{2} \partial_x \overline{(\psi'^2_x - \psi'^2_y)} + \partial_y \overline{\psi'_x \psi'_y} + \overline{h\psi'_x} \right] \approx -\beta^{-1} \overline{h\psi'_x} \end{aligned}$$

Eddy fluxes

$$D = \frac{d}{dt} \frac{1}{2} \overline{\eta^2} = \beta^{-2} \frac{d}{dt} \frac{1}{2} \overline{q'^2} = \beta^{-1} \overline{v'q'} = -\beta^{-1} \overline{(\nabla^2 \psi' + h)\psi'_x}$$

$$\approx -\beta^{-1} \left[\frac{1}{2} \partial_x (\overline{\psi_x'^2} - \overline{\psi_y'^2}) + \partial_y \overline{\psi'_x \psi'_y} + \overline{h\psi'_x} \right] \approx -\beta^{-1} \overline{h\psi'_x}$$

change to spatial average

scale separation

$$P^{(x)} = \overline{h\psi'_x} \approx -\overline{\psi' h_x}$$

Topographic
form stress

Eddy fluxes

$$D = \frac{d}{dt} \frac{1}{2} \overline{\eta^2} = \beta^{-2} \frac{d}{dt} \frac{1}{2} \overline{q'^2} = \beta^{-1} \overline{v'q'} = -\beta^{-1} \overline{(\nabla^2 \psi' + h) \psi'_x}$$

$$\approx -\beta^{-1} \left[\frac{1}{2} \partial_x (\overline{\psi_x'^2} - \overline{\psi_y'^2}) + \partial_y \overline{\psi'_x \psi'_y} + \overline{h \psi'_x} \right] \approx -\beta^{-1} \overline{h \psi'_x}$$

change to spatial average

scale separation

$$P^{(x)} = \overline{h \psi'_x} \approx -\overline{\psi' h_x}$$

Topographic
form stress

Discuss: really?

Summary

- An initially turbulent flow above topography tends to a minimum enstrophy steady solution that is approximately along isobaths.
- On a beta-plane, the minimum enstrophy solution implies an westward interior flow.
- Topographic form stress appears to play a key role in driving the large-scale flow.

Discussion topics

(Or things to think about in the privacy of your own study)

- “The preferred topography”: $|\hat{h}|^2(|k|) \propto |k|^{-2}$
- What is the relevance of freely-decaying solutions to the understanding of real (forced) geophysical flows?
- Differences between closed-basin and doubly-periodic solutions.
- The $h \rightarrow 0$ limit.

The limit $h \rightarrow 0$ is a bit murky...

A simple eigenproblem

$$\nabla^2 \psi = \mu \psi$$

The minimum enstrophy solution is

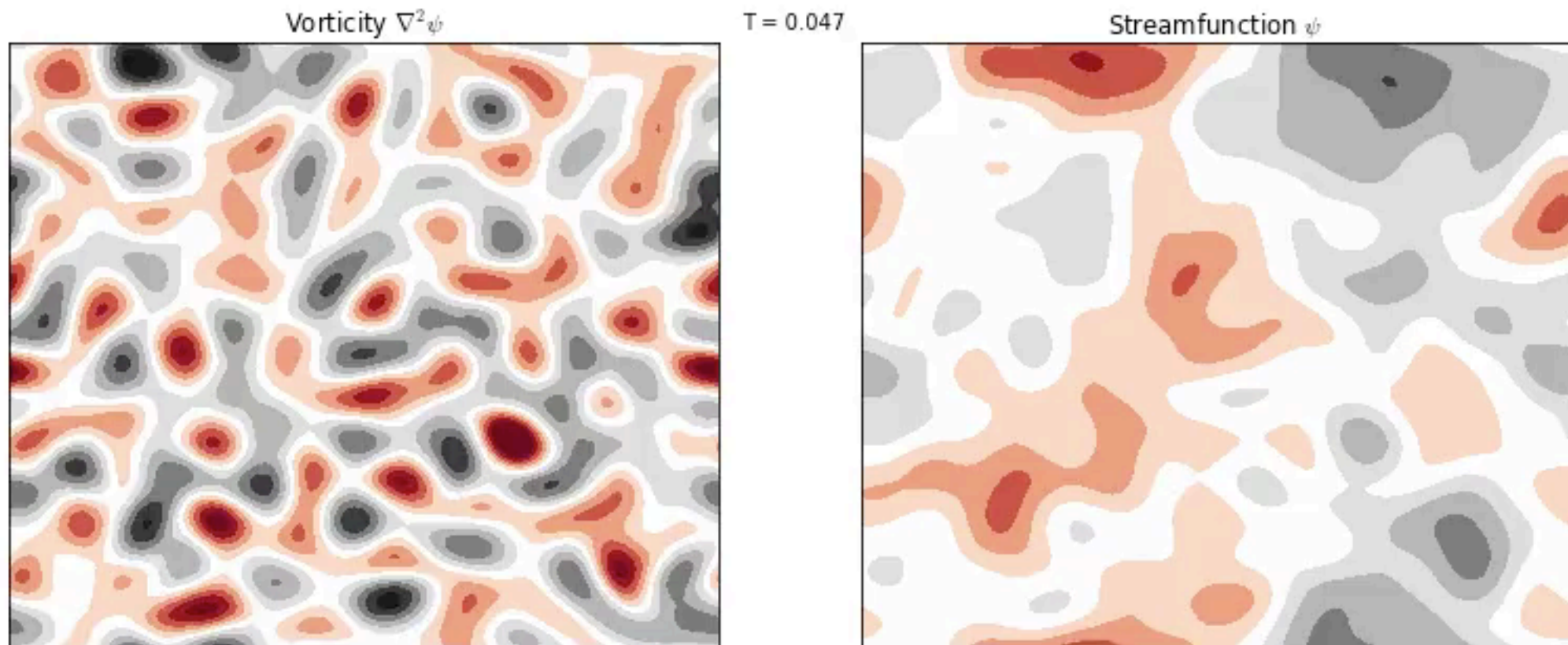
$$\psi_0 = A \cos kx + B \cos ly$$

$$k^2 + l^2 = 1 \quad \rightarrow \quad \mu = -1$$

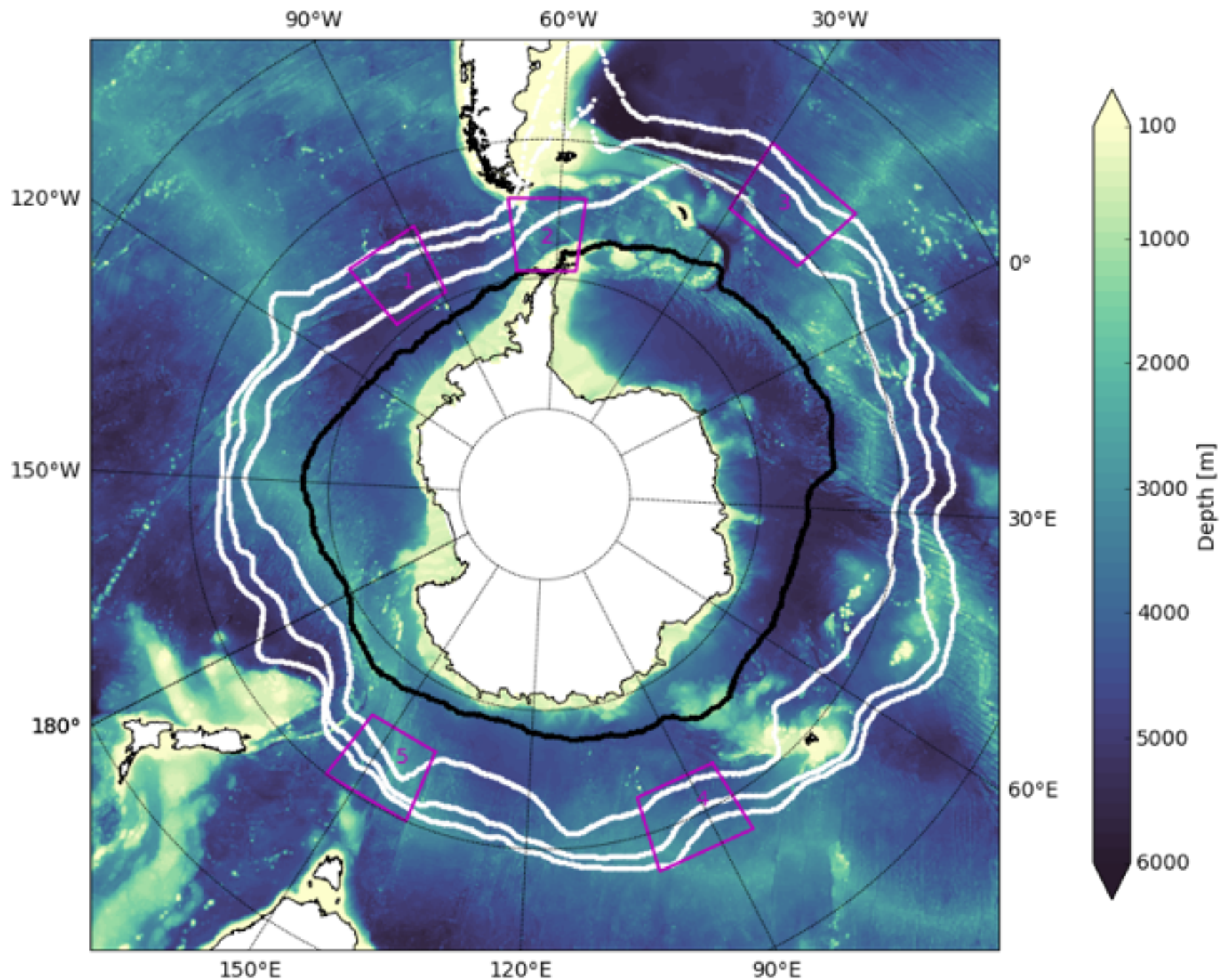
Discuss: What determines A and B?

2D turbulence

An initial value problem ($h=0$)

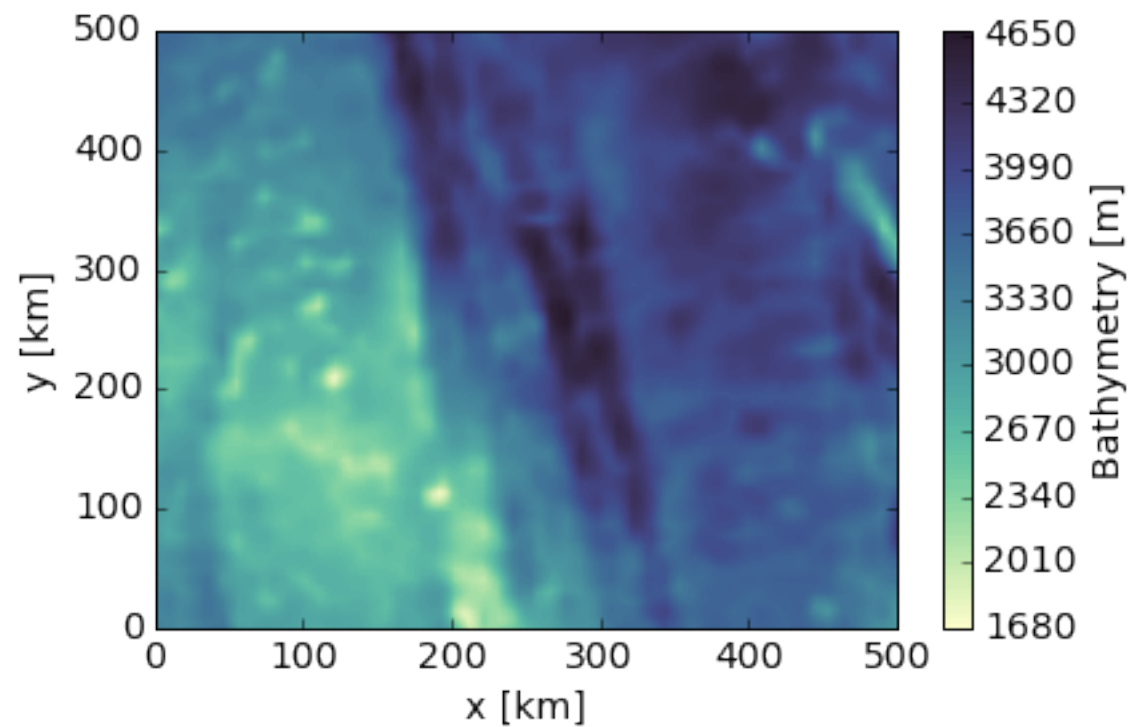


Southern Ocean Topography

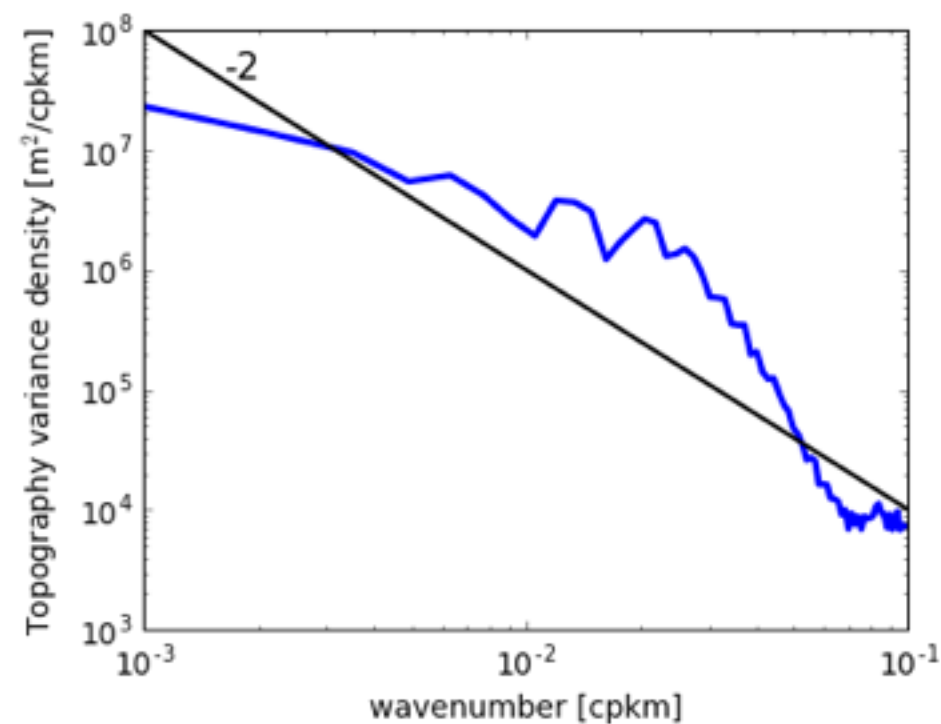
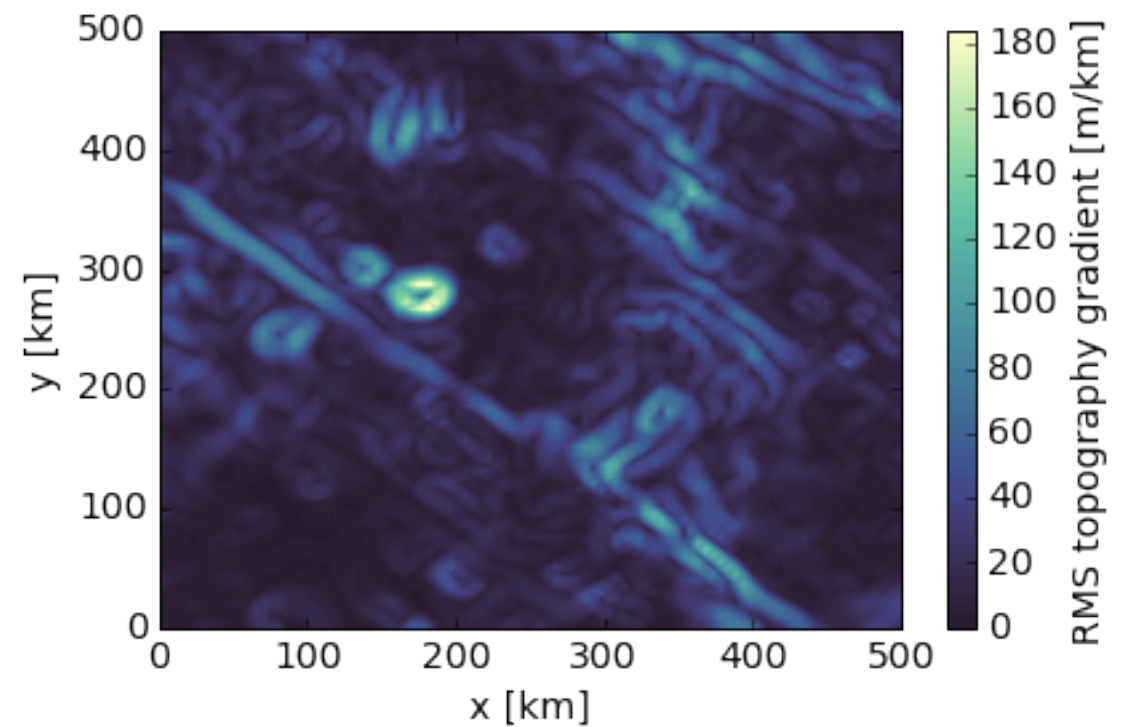


Drake Passage Topography

Topography h



Topographic gradient



$$K = \beta_{topo} / \beta = \frac{\frac{|f_0|}{H} |\nabla h|_{rms}}{\beta} \approx 54.24$$

$$H = \langle h \rangle \approx 3716 \text{ m}$$

$\langle \cdot \rangle$ spatial average